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TECHNICAL NOTE

FLUTTER INVESTIGATION OF A TRUE-SPEED DYNAMIC MODEL WITH

VARIOUS TIP-TANK CONFIGURATIONS

By John L. Sewall, Robert W. Herr, and William B. Igoe

Langley Research Center Langley Field, Va.

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FLUTTER INVESTIGATION OF A TRUE-SPEED DYNAMIC MODEL WITH

VARIOUS TIP-TANK CONFIGURATIONS*

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SUMMARY

A 1/6-scale dynamic model, equipped with wing tip tanks and representative of unswept fighter-type airplanes of conventional plan form, was tested for flutter. The model was dynamically scaled to flutter at the same speed as a full-scale configuration. The parameters important to flutter were satisfactorily approximated by the use of a spar-balsa segment-type wing construction.

A flutter-arresting device was located within each tip tank which would provide for a sudden shift in the tip-tank center of gravity. This device was very effective in stopping symmetric flutter for cases where the tip-tank center of gravity was shifted forward of the elastic axis of the wing.

The flutter investigation was concerned with the effects of external stores whose inertial and geometric properties were systematically varied. In general, the results obtained concur with those found in previous researches in that symmetric flutter speeds were increased for tip-tank centers of gravity forward of the wing elastic axis, and tended to decrease as the ratio of uncoupled wing bending frequency to torsion frequency approached unity. Increasing the volume of the tip tanks by a factor of 3 tended to decrease the flutter speed. The addition of a horizontal fin at the rear end of each tank may have been beneficial for the smaller tank and slightly detrimental for the larger tank.

Flutter speeds were calculated by means of a conventional Rayleigh-Ritz type of flutter analysis employing uncoupled modes. The effects of various assumptions in the calculations involving structural damping, higher structural modes, compressible aerodynamic coefficients, and sting-mount flexibility were explored. The calculations were, in general, excessively conservative for tip-tank centers of gravity near the elastic axis; however, the introduction of structural damping materially improved the agreement between experimental and calculated flutter speeds for ratios of uncoupled wing bending frequency to torsion frequency near unity. The introduction of sting-mount flexibility showed that this could have a

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strong effect on flutter speed particularly in those cases where the natural frequency of the sting system was near the flutter frequency, and further work in this regard is warranted in future flutter studies of sting-mounted models. Also indicated as desirable is the need for continued research on the oscillating aerodynamic coefficients for flutter calculations on wings equipped with large external stores or nacelles.

INTRODUCTION

Within the past several years a number of cases of flutter are known to have occurred on fighter and light bomber airplanes (at least in their development stages) carrying large external stores. Considerable interest has, therefore, developed in flutter problems for wings carrying such stores. References 1 to 7 are representative of both the theoretical and experimental investigations of problems of this kind. The calculations of flutter instabilities for these cases have, in general, been unsatisfactory where two-dimensional aerodynamic forces and moments together with a limited number of structural modes have been employed, and where aerodynamic forces and moments on the tip tank have been ignored. Because of weaknesses inherent in many of the more complete analyses available at present, greater emphasis has come to be placed on wind-tunnel testing of dynamically scaled models with external stores.

Although valuable experience has been obtained from flutter tests with scaled speed models (see, for example, refs. 8 to 10), these tests cannot provide needed information on the effects of Mach number. The extrapolations from low-speed model results to full-scale high Mach number conditions have proved troublesome and ambiguous, particularly where the margins of safety from flutter may have been small. In order to circumvent this difficulty, it is necessary to scale the model so that its Mach number at flutter will be the same as that of the prototype, or, in other words, to design and test a true-speed model. Accordingly, two principal objectives of the investigation reported in this paper were (1) to determine the feasibility of designing and building a true-speed dynamically similar model with tip tanks to meet a given set of elastic and inertial specifications, and (2) to locate the experimental flutter boundaries of this model for a variety of tip-tank conditions. Another objective was to evaluate a flutter-arresting device involving a rapid shift of the tip-tank center of gravity and having application to fullscale-flight flutter-testing techniques.

For the design of the model, in addition to keeping its flutter speed the same as that of the prototype, the principles followed in the scaling required that nondimensional parameters important in flutter, such as mass ratio, frequency ratio, and reduced frequency, remain the same on the model as on the full-scale counterpart. No attempt was made

to scale the stiffnesses in the tip-tank or wing-tip-tank attachment. Recent investigations reported in references 11 and 12 have dealt with the inertial and damping properties of sloshing fuel; however, in the present investigation fuel was simulated by means of solid weights, and no effort was made to account for the effects of fuel sloshing.

The main effort of the wind-tunnel test program was directed toward location of the flutter boundaries for a wide range of tip-tank inertial conditions corresponding to ratios of uncoupled wing bending frequency to torsion frequency near unity. Variations in these tip-tank conditions were accomplished by the use of fixed, solid weights. Also investigated were the aerodynamic effects of tip-tank size and the effects of a horizontal fin mounted at the rear portion of each tip tank. In order to study asymmetric flutter, the model was provided with one degree of freedom in roll.

The results of the flutter experiments are compared with the results of flutter calculations performed by utilizing a Rayleigh-Ritz type of analysis employing two-dimensional aerodynamic forces and moments together with a combination of uncoupled vibration modes to represent the flutter mode. No aerodynamic forces and moments were assumed to act on the tip tank.

The scaled wing stiffnesses and inertial properties specified in the design of the model were chosen to be fairly representative of present-day fighter-type airplanes. Scaled wing stiffnesses and inertial properties of an actual fighter airplane are included as a comparison. Also included for purposes of comparison are the results of a flight flutter experience on this airplane involving special tip-tank inertial conditions.

SYMBOLS

- a nondimensional wing elastic axis location relative to halfchord, positive for elastic axis rearward of midchord
- b half-chord of wing, ft
- b_F flange width on wing spar, in. (see fig. 9)
- br reference half-chord of wing (taken at the station at threefourths of the semispan), ft
- b_W web spacing on wing spar, in. (see fig. 9)

 g_{α_1}

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root chord of wing along fuselage center line, in. (see fig. 3(a))
c_{\mathbf{r}}
           tip chord of wing at y = l_w = 36, in. (see fig. 3(a))
c_{t}
           distance between wing elastic axis and center of gravity, in.;
е
              positive for center of gravity rearward
           distance between wing elastic axis and tip-tank center of
e<sub>t.</sub>
              gravity, in.; positive for center of gravity rearward
f
           frequency, ω/2π, cps
f_{b_1}
           natural first bending frequency (experimental), cps
           natural second bending frequency (experimental), cps
fb2
\mathbf{f}_{\mathbf{f}}
           flutter frequency, cps
f_{h_l}
           uncoupled first bending frequency (calculated), cps
\mathbf{f}_{h_2}
           uncoupled second bending frequency (calculated), cps
           uncoupled frequency of effective sting-fuselage combination,
f_s
              cps (see appendix)
f_{t_1}
           natural first torsional frequency (experimental), cps
^f\alpha_l
           uncoupled first torsional frequency (calculated), cps
\mathbf{f}_{\alpha_{\!\scriptscriptstyle \mathcal{D}}}
           uncoupled second torsional frequency (calculated), cps
           coefficient of structural damping based on logarithmic decrement
g
              g = \frac{1}{n\pi} \log_e \frac{a_0}{a_n} where a_0 is the initial amplitude and a_n
               the nth amplitude in n cycles of freely decaying oscillation
h_{\mathbf{w}}
           flange spacing on wing spar, in. (see fig. 9)
           reduced frequency, bw/v
k
           coefficient of structural damping in first bending mode
g_{h_1}
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coefficient of structural damping in first torsional mode

- k_s spring constant of equivalent sting support system, lb/in. (see appendix) 2 general dimension of length (see text on selection of scale factors) ฆ wing semispan from center line of fuselage to center of wingtip-tank attachments, in. mass of wing per unit length, slugs/ft m distance from effective pitch axis of rotation of sting to wing r_s elastic axis at fuselage center line (see appendix) \mathbf{r}_{α} nondimensional radius of gyration of wing section about elastic axis, $/I_{\alpha}/mb^2$ flange thickness on wing spar, in. (see fig. 9) tϝ web thickness on wing spar, in. (see fig. 9) tw velocity, fps flutter speed, fps v_f W weight of wing per unit length, lb/in. \mathbf{x}_{α} nondimensional center-of-gravity location of wing section relative to half-chord, e/12b; positive for center of gravity rearward of wing elastic axis У spanwise coordinate of wing as measured from fuselage center line, in. wing bending stiffness, lb-in.2 EI F + iGoscillating aerodynamic functions for two-dimensional incom-
- pressible flow (see ref. 13)

 GJ wing torsional stiffness, lb-in.²
- I_{f_r} mass rolling moment of inertia of one-half fuselage about fuselage center line, ft-lb-sec²
- I_{f_S} mass pitching moment of inertia of one-half fuselage and effective sting combination about effective axis of rotation in sting, ft-lb-sec² (see appendix)

 $\mathbf{I}_{\mathrm{t_r}}$ mass rolling moment of inertia of tip tank about wing tip, $\mathrm{ft\text{-}lb\text{-}sec^2}$

 I_{α} mass pitching moment of inertia of wing per unit length about wing elastic axis, lb-sec²

 $I_{\alpha t}$ mass pitching moment of inertia of tip tank about wing elastic axis, in-lb-sec 2

M Mach number

R Reynolds number

W weight, lb

Wt weight of tip tank, lb

Ww weight of each wing panel, 1b

 η_{asym} nondimensional span coordinate used in antisymmetric flutter calculations, y/l_W

 η_{sym} nondimensional span coordinate used in symmetric flutter calculations, $\frac{y-3}{h_{\text{tr}}-3}$

 κ wing mass ratio, $\pi \rho b^2/m$

 λ scale factor for length, $l_{\rm M}/l_{\rm F}$ or $b_{\rm M}/b_{\rm F}$

 μ tip-tank weight ratio, W_t/W_w

ρ density of test medium, slugs/cu ft

ω angular frequency, 2πf, radians/sec

Subscripts:

f flutter except as noted

h bending degree of freedom

r reference station at three-fourths semispan

t tip tank

w wing

f full-scale airplane except as noted

M model

S sting support system

W spar web

α torsional degree of freedom

asym antisymmetric fuselage boundary condition (fuselage free to roll) considered in flutter calculations

sym symmetric fuselage boundary condition (cantilever and fuselage pitch) considered in flutter calculations

DESCRIPTION OF MODEL

The model (see fig. 1) was sting-mounted in the Langley 16-foot transonic tunnel (fig. 1(a)) and was representative of unswept-wing fighter-type airplanes equipped with wing-tip fuel tanks. The aspect ratio of the wings without tip tanks was 6.0, the taper ratio was 0.381, and the airfoil shape was an NACA 65AO13 section. The wing was untwisted, and the 52-percent-chord line was normal to the longitudinal center line of the airplane; although the elastic axis was swept back $3\frac{10}{2}$, the wing was considered unswept.

The model was tested with two different-size tip tanks (shown in fig. 2) which were mounted so that the longitudinal axis of the tank fell in the plane of the wing. Tests were conducted with and without horizontal fins located at the rear of the tip tank as shown in figure 2. Some wing details may be seen in figure 3. Each wing panel was mounted as a cantilever (see fig. 3(a)) to a cylindrical body simulating the fuselage as a rigid structure that provided freedom of the model in roll.

Selection of Scale Factors

Scaling of the model was based on the simplified model construction discussed in reference 14 (ch. 11). The quantities scaled were linear dimensions, mass, mass unbalance, moments of inertia, frequencies, and stiffnesses of the prototype. The wind tunnel used had an octagonal, slotted throat and was capable of attaining Mach numbers up to 1.05.

The following table lists the flutter parameters considered significant in this investigation in terms of the geometric scale factor λ . The model was chosen to be 1/6 the scale of a representative fighter-type airplane, and, therefore,

$$\lambda = \frac{b_{\underline{M}}}{b_{\underline{F}}} = \frac{1}{6}$$
 or $\lambda = \frac{l_{\underline{M}}}{l_{\underline{F}}} = \frac{1}{6}$

where the subscripts M and F refer to model and full scale, respectively. The choice of this particular value of λ was influenced by the desire to obtain Reynolds numbers that would be as realistic as possible with respect to the full-scale airplane within the restriction imposed by available wind-tunnel size. The density factor ρ_{M}/ρ_{F} was assumed to be unity.

Parameter	Symbol	Scale factor
Mass per unit length	$\frac{\frac{m_{\mathbf{M}}}{m_{\mathbf{F}}}}$	$\lambda^2 = \frac{1}{36}$
Mass	$rac{m_{ m M} l_{ m M}}{m_{ m F} l_{ m F}}$	$\lambda^3 = \frac{1}{216}$
Mass moment of inertia per unit length	$rac{ extsf{I}_{lpha_{ extsf{M}}}}{ extsf{I}_{lpha_{ extsf{F}}}}$	$\lambda^{14} = \frac{1}{1296}$
Mass moment of inertia	$\frac{\mathtt{I}_{\alpha_{\!\scriptscriptstyle \mathbf{M}}} \imath_{\!\scriptscriptstyle \mathbf{M}}}{\mathtt{I}_{\alpha_{\!\scriptscriptstyle \mathbf{F}}} \imath_{\!\scriptscriptstyle \mathbf{F}}}$	$\lambda^5 = \frac{1}{7776}$
Frequency	$\frac{\omega_{\mathrm{M}}}{\omega_{\mathrm{F}}}$	$\frac{1}{\lambda} = 6$
Bending stiffness	(EI) _M (EI) _F	$\lambda^{4} = \frac{1}{1296}$
Torsional stiffness	$\frac{(GJ)_{M}}{(GJ)_{F}}$	$\lambda^{4} = \frac{1}{1296}$
Structural damping	$\frac{\left(g_{h}\right)_{M}}{\left(g_{h}\right)_{F}}$ or $\frac{\left(g_{\alpha}\right)_{M}}{\left(g_{\alpha}\right)_{F}}$	1
Mach number	M _M M _F	1

L δ 1 It should be noted that $M_M = M_F$ is not precisely true because of differences which may exist between the stagnation temperature in the tunnel and that in free air for an airplane. No attempt was made to design a given value of structural damping (g_h,g_α) into the model. Accurate measurements of the bending and torsional damping coefficients, as determined from logarithmic decrements during the test program, were difficult. The coefficients measured varied from 0.01 to 0.10. The mean value appeared to be approximately 0.035 for both bending and torsion degrees of freedom. This value agrees well with a structural damping value of 0.044 obtained from a comparable full-scale fighter airplane.

Model-Construction Details

<u>Wings</u>.- The solid curves shown in figures 4 to 8 give the scaled stiffness and inertial properties typical of current fighter-type airplanes, and the model was designed to meet these specifications.

The wings were of spar-segment construction as shown in figure 3. The spar, the main structural member, was essentially a built-up box section made from thin sheet steel and welded at all four corners as shown in section A-A of figure 3(a). This type of spar construction was chosen because of the relative ease with which the bending and torsional stiffnesses could be varied essentially independent of one another. Dimensions of the spar are given in figure 9.

The wing plan form and airfoil shape were obtained by assembling a series of adjacent wing segments to the spar in the manner shown in figure 3(b) for a typical segment. Ballast weights of the kind indicated in the figure were added to each segment to adjust the wing mass, mass unbalance, and mass moment of inertia to the design values. The segments were built-up hollow sections made from sheet balsa and were coated inside and out with fiber glass and Paraplex to provide extra rigidity for the segment to withstand better the forces encountered during violent flutter. Each segment was fastened to the spar flanges by means of four screws, each screw bearing on a small metal-backed strip of hardwood glued to the inside of the segment.

The measured inertial properties of the completed model are indicated in figures 4, 5, and 6, and the spanwise distributions of wing stiffnesses in bending and torsion are given in figures 7 and 8, respectively. The attachment of the balsa segments to the spar caused no measurable increases in bending and torsional stiffnesses. The gaps between the balsa segments were covered with stiffened paper for most of the flutter runs, and later this material was replaced with silk for the remaining runs. Although the stiffness contribution of the silk was more noticeable than that of

the paper, the measured stiffnesses of the assembled model were not significantly altered because of the use of either material. However, the measured stiffnesses which are presented in figures 7 and 8 were obtained prior to the flutter tests without either paper or silk seals attached.

Stiffness measurements made at different times during the construction of the model and the flutter test program revealed some weakening of the wing due to flutter violence; however, on the basis of frequency checks made during the flutter program, it appeared that such changes had little effect on the vibration characteristics of the wing. The dashed-line curves in figures 7 and 8 were considered most representative of the stiffness properties of the wing and were used in the flutter calculations.

Scaled stiffness and inertia parameters of an actual fighter airplane are included in figures 4 to 8 for comparison purposes. For the inertial parameters (figs. 4 to 6), these curves apply to full fuel load in the integral wing fuel tanks.

Tip tanks. Figure 2 shows the two tip tanks that were used in this investigation. The smaller of these two tanks (hereinafter referred to as tank A) was a geometrically scaled reproduction of a full-sized 230-gallon tip tank, whereas the larger tank (tank B) was designed to have the same shape but three times the internal volume of the smaller tip tank. No attempt was made to scale the stiffness of the tip tank or its attachment to the wing.

Construction of the model tip tanks consisted of a duralumin cylindrical center section with hollow plastic nose and tail cones made of fiber glass and Paraplex. Tank B was assembled by replacing the smaller nose and tail cones with larger ones and mounting a thin duralumin cover (shown removed in fig. 2) over the center section. For each tank, two tail cones were provided, one with a horizontal fin and one without a fin, in order to explore the effects of the fin on flutter. For tank A the fin was scaled geometrically from the full-sized configuration. Each tip tank housed a flutter-arresting device which is described in detail later in the paper. This device involved the use of one of three different shifting weights to achieve a rapid change in the mass unbalance of the tip tank. Table I gives the inertial properties of the tank for each of these weight conditions with no additional weights added.

The geometrical properties of both tip tanks and fins are also given in table I. The fairing at the juncture of the wing and tank A can be seen in figure l(b). This fairing was completely covered by the center section of tank B.

Figure 10 indicates the manner in which additional weights were mounted in the tanks to simulate a given fuel load. As previously noted, no effort was made to take into account the effective changes in the tiptank inertial characteristics due to fuel sloshing. The tip-tank weight is hereinafter referred to in terms of the weight ratio μ which is defined as the ratio of the tip-tank weight to wing-panel weight by the relation

$$\mu = \frac{W_{t}}{W_{w}} = \frac{W_{t}}{10}$$

since $W_w = 10$ pounds.

Figure 11 is included for the purpose of relating μ to full-scale fuel loads for both size tip tanks. In view of the results of reference 11 concerning the effective weight of sloshing fluid in a tank, the fuel loads simulated in the present investigation by means of solid weights correspond to increased actual fuel loads. Although tank B had three times the internal volume of tank A, the full-scale fuel capacity of both tanks was considered to be the same. The empty weight of tank B was based on a further assumption that the larger full-scale tank housed fixed equipment whose weight was approximately 50 percent of the weight of its full fuel load. The weight of this fixed equipment accounts for the large differences in μ corresponding to the tank-empty condition in figure 11. The minimum weight condition indicated by the dotted boundary in the figure, and applying to a solid fuel load of 25 percent full, corresponds to the lightest weight of the flutter-arresting device for tank A with fin. (See table I.)

Fuselage .- The main structural member of the fuselage was a hollow steel section to which the wings were mounted through steel brackets welded to the root of each wing spar. This center section was made hollow so as to permit passage of the instrument wires leading from the model to the control room. The fairings forming the external contours of the fuselage were attached to either end of the main member, as may be seen in figure 1(b) which shows model with middle fairings removed. The elastic properties of the fuselage were not intended to be representative of current fighter designs; the fuselage was, in fact, essentially rigid with respect to the wings. Bearings were housed at either end of the center section to permit a rolling degree of freedom of ±180°. Low stiffness springs were housed within the fuselage to position the model horizontally in the tunnel when the model was given freedom to roll. Small aerodynamic control tabs (fig. 1(a)) were located at the rear end of the fuselage to provide manual control of the model in roll. The rolling degree of freedom could be completely locked out so as to obtain essentially the cantilever or symmetric boundary condition. The actual measured rolling inertia of the model fuselage was 2.36 in-lb-sec² which is somewhat larger than the scaled value of a typical fighter fuselage.

Instrumentation.— All the frequencies in both the natural vibration survey and the flutter investigation were measured by means of resistance wire strain gages located at three different stations along each wing spar as indicated in figure 3(a). The bending gages were mounted to the inside surfaces of the overhanging flanges, and the torsion gages were mounted to the vertical web members. The signals from these strain gages were fed into a recording oscillograph. The position of the model in roll could be determined by means of a slide-wire position indicator located at the forward end of the main center-section member of the fuselage.

In order to determine the vertical acceleration of the fuselage during flutter, an accelerometer was located on the main center-section member at the juncture of the fuselage center line and the wing elastic axis.

High-speed motion pictures were taken of the model during flutter from a position outside the tunnel and in line with the span of the wing. The only other item of model instrumentation was associated with the operation of the flutter-arresting device, discussed in the following section.

Flutter-arresting gear .- A desirable and important feature of windtunnel flutter testing is a reliable means of preventing a destructive buildup of amplitude once a flutter condition is attained. The more common methods used at low speeds such as restraining wires or mechanical means within the flow are not too satisfactory at transonic speeds. Several alternative methods have been suggested. The Wright Air Development Center (WADC), for instance, has successfully restrained rudder flutter in flight by quickly changing the mass balance of the control surface. One of the objectives in this series of tests was to evaluate a flutter-arresting device involving a rapid shift of the tip-tank center of gravity and its possible application to flight flutter-testing techniques. The shift of the tip-tank center of gravity was accomplished by moving a piston, within a cylinder running the length of each tank, from its extreme rearward position to its extreme forward position or conversely. The piston could not be stopped at an intermediate position. The position of the piston was controlled by an observer or by an electronic device that monitored the electrical outputs of one of the strain gages on the wings. When the oscillating stress at that station exceeded a predetermined level, the piston was automatically fired to its other position.

Included in figure 3(a) is a simplified sketch of the flutterarresting gear. Forward travel of the piston was obtained by supplying The mass of the cylinder, being only 0.53 pound or 6 percent of the tank A full weight, imposed no weight penalties except in cases where low fuel loads were to be simulated. The pistons used weighed 0.40, 0.78, and 1.07 pounds, corresponding respectively to 4.8, 9.2, and 12.7 percent of the tank A full weight.

MODEL TESTS

Vibration Survey

Before the flutter tests were started, natural frequencies of vibration of the model were measured with the model mounted in a horizontal position on a dummy sting. The flexibility of this sting was made approximately equivalent to that of the sting in the Langley 16-foot transonic tunnel so that the model frequencies obtained in this survey would be comparable with those obtained on the sting in the tunnel. The model was excited both manually and with an electrodynamic shaker attached at various points along the wing.

Wind-Tunnel Test Procedure

Figure 1(a) shows the model mounted on the sting in the Langley 16-foot transonic tunnel. The flutter tests were conducted at Mach numbers up to 0.94 and at Reynolds numbers per foot up to 4.2×10^6 .

In conducting the flutter tests, the tunnel speed was increased gradually and maintained at various speeds while the model response to random tunnel disturbances was observed and recorded with the flutterarresting piston in the forward and rearward positions of the tip tank. Since only two positions of the piston were possible (extreme forward or extreme rearward) during the approach to flutter, the piston was left in the position that appeared to have the least damping as observed by the aforementioned response of the model to the random tunnel disturbances.

The wind-tunnel tests were terminated either when flutter was encountered or when the test could not proceed further, because of tunnel power limitations or lack of effectiveness of the flutter-arresting device.

When flutter was encountered, the tunnel conditions were recorded simultaneously with the operation of the recording oscillograph and movie camera. At the same time, the piston was fired to the position which appeared to have more damping, and it was hoped that this would result in a condition that was not in a flutter region (a condition which was usually obtained). The oscillograph records and movies were generally taken continuously before, during, and immediately after the firing of the piston.

The model was excited manually by plucking the wing tips in bending and torsion prior to and often after each flutter run, and records were made of the bending and torsional vibration frequencies. These frequencies were measured after the flutter runs to provide checks on the structural integrity of the model.

RESULTS AND DISCUSSION

For convenience most of the experimental and calculated results of this investigation are indexed in table II. Figures 12 to 18 listed in this table show flutter Mach number plotted against $e_{\rm t}$ for both experimental and calculated results. Figures 19 to 23 listed in the table show frequency plotted against $e_{\rm t}$.

Experimental Results

Table III compares some of the results of the vibration survey with the results of vibration tests on an actual fighter-type airplane with tip tanks. The comparisons are made between model frequencies obtained during the vibration survey and scaled frequencies for the closest wing and tip-tank inertial conditions available for both model and airplane. The airplane frequencies and tip-tank inertial conditions were scaled according to the scale factors listed earlier in the section entitled "Description of Model." The agreement shown by the limited comparisons in table III for first bending and first torsion indicate how well the flutter model may have represented the full-scale configuration. The fact that the two scaled second bending frequencies listed for the airplane were considerably higher than the model second bending frequencies could be partly due to reduced fuel rolling inertia effects (see ref. 12) which may have been present in the airplane but were not realized in the model.

Table IV presents the results of the flutter tests in which the tiptank inertial properties were systematically varied so as to obtain flutter for ratios of uncoupled wing bending frequency to torsion frequency f_{h_1}/f_{α_1} near unity. In this table the tip-tank weight ratio μ and moment of inertia for each flutter run are represented by a simple designation. In this designation the letter "A" or "B" identifies the tip tank used, the set of numbers following the first dash refers to μ , and the set of numbers following the second dash gives the uncoupled frequency ratio f_{h_1}/f_{α_1} . As previously noted, the value of μ was determined from the ratio of tip-tank weight to wing-panel weight; the frequency ratio was determined from figures 24 and 25. The phase angle indicated in table TV is the angle by which the bending strain led the torsion strain.

Effect of tip-tank center of gravity.— The effect of e_t on flutter speed may be seen in figures 12 to 18 for various values of μ and frequency ratio f_{h_1}/f_{α_1} . The various frequency ratios correspond to different moments of inertia for a given value of μ . Speeds where flutter was encountered are represented by the solid points and the open points indicate speeds reached in the tests without encountering flutter.

With the exception of figure 16, the results presented in figures 12 to 18 were obtained with the fuselage locked to prevent roll; thus, these results are for symmetric flutter, that is, flutter with both wing panels oscillating in phase. These figures, except for figure 12, apply to the smaller tip tank (tank A). In general, as may be observed, flutter speeds are increased for center-of-gravity positions forward of the elastic axis and tend to decrease as the ratio of uncoupled wing bending frequency to torsion frequency approaches unity. (See, for example, fig. 14.)

The flutter frequencies corresponding to the flutter speeds (shown in figs. 12 to 16) are presented in figures 19 to 23. In addition, figures 19 to 23 show the first and second measured wing natural frequencies with the model on the sting, the first and second calculated coupled frequencies, and the first bending and first torsional uncoupled frequencies. Each figure applies to a given tip-tank weight ratio μ and a given calculated frequency ratio $f_{h_1}/f_{\alpha_1}.$ The fact that the experimental flutter frequency consistently fell between the measured first bending and first torsional wing frequencies of the model mounted on the sting indicates that the wing with tip tanks fluttered mainly in a combination of the first bending and first torsional structural modes. This type of motion was observed on the model during flutter. To show this, a typical flutter cycle, figure 26, is reproduced from the motion pictures taken during flutter.

Effect of tip-tank weight.— A comparison of the data of figure 13 $\left(f_{h_1}/f_{\alpha_1}=0.93;\ \mu=0.59\right)$ with the data of figure 15 $\left(f_{h_1}/f_{\alpha_1}=0.92;\ \mu=1.005\right)$ indicates that an increase of μ from 0.59 to 1.005 is accompanied by only a moderate decrease in flutter speeds for any given value of e_t . However, the flutter data given in table IV are insufficient to show the effect at low values of μ .

An appraisal of the flutter-arresting gear.— In figures 12 to 15 the effect of the tip-tank center of gravity on the symmetric flutter boundary indicates that, in general, the forward position of the piston should have resulted in a flutter-free condition. The fact that this proved to be the case is attested to by the motion pictures and the oscillograph records which were taken continuously when flutter occurred to provide complete time histories of the flutter response and subsequent damping action caused by the change in position of the piston. Figure 27 is a typical flutter oscillogram which shows clearly how the forward position of the piston damped the oscillations. For all conditions where the tip-tank center of gravity was shifted forward of the elastic axis, the flutter-arresting gear was effective in damping flutter.

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Despite the effectiveness of this device, the shapes of the flutter boundaries were such as to impose certain restrictions on the practical application of the variable unbalance feature of the device. These restrictions are evident in a comparison of figures 15 and 16 which show increasing flutter speeds for tip-tank centers of gravity rearward of certain locations. Figure 15 shows that it was possible to traverse a small portion of the flutter region into a safe region without dangerous amplitudes building up. However, figure 16 shows a flutter condition where this was not possible. With the center of gravity approximately 3.6 inches rearward of the elastic axis, firing the piston forward established a condition also very close to the flutter boundary as indicated by the open point at approximately 1.4 inches rearward of the elastic axis. Since this shift did not change the center of gravity to a position in the safe region, this test could not be safely conducted beyond a Mach number of 0.37.

A further restriction in the application of the arresting device to flight-flutter use may be seen by examining the weight and size of a full-scale flutter-arresting device. If a full-scale device of this type were reproduced according to the scale factors applicable to the flutter model, the flutter arresting tube would be 6 times as long, or approximately 12 feet. The full-scale piston would be 216 times as heavy, or would weigh approximately 230 pounds for the heaviest piston.

The cumbersome features of a mass of this size moving such a distance in a full-scale tank may be offset by the time of travel of the piston from one end of the tube to the other. In order for the piston travel

to be accomplished in the same number of flutter cycles for both model and airplane, the full-scale piston would have to traverse the arresting tube length in about $2\frac{1}{2}$ seconds.

It is recognized that flutter conditions can occur which may be insensitive to a change in mass unbalance and that under such conditions more suitable flutter-arresting devices may be those which suddenly cause a large increase in the damping of the structure. However, for the symmetric flutter encountered in the present study, the shifting mass-unbalance principle proved very effective within the limitations discussed in the preceding paragraphs.

Effect of tip-tank size. The effect of the size of the tip tank on flutter can be seen in figure 12 by a comparison of the curve through the large symbols (pertaining to the large tank) with the curve through the small symbols (pertaining to the small tank), both tanks having nearly the same value of μ and moment of inertia (and, as a result, nearly the same frequency ratio f_{h_1}/f_{α_1}). The reduction in flutter speed for the large tip tank, which was as much as 15 percent in the region covered by the tests, is probably due mainly to the increased aerodynamic forces resulting from the larger tank, although it should be noted that approximately $l\frac{1}{2}$ percent of the area of the half-wing was covered by the large tip tank. (See fig. 3(a).)

Effect of tip-tank fins. - Most of the flutter runs were conducted with a horizontal fin attached near the trailing edge of the tip tank on the outboard side. Some tests were conducted with the fin removed, and the effect of this alteration on the flutter speed may be seen by comparing certain tests in table IV for both large and small tip tanks. In some of these cases the fin appeared to be beneficial and in other cases it was not. For tip tank A, for example, comparisons between tests 4 and 5 for a low tip-tank weight and between tests 25 and 26 for a high tip-tank weight show that the use of the fin significantly increased the flutter speed. In contrast, however, comparisons between tests 6 and 7 for tip tank A and between tests 40 and 41 for tank B indicate that the addition of fins had little effect on the flutter speed and actually reduced the flutter speed slightly for the large tank.

Observations on fuselage freedoms.— The effect of freeing the fuselage to roll about its longitudinal axis is shown in figure 16. These data pertain to the large tip tank for $\mu=1.02$ and an uncoupled symmetric frequency ratio of 1.04. The solid curve through the circles corresponds to the fuselage-locked condition, whereas the squares apply to tests in which the rolling freedom was permitted. This additional freedom produced little or no effect on the flutter speed in the region where

symmetric flutter was encountered. Here, as in previous figures, the open points represent speeds reached without encountering flutter. Previous experiences with antisymmetric flutter (see, for example, ref. 10) have indicated that this type of flutter can be critical for forward tiptank centers of gravity. In that case, firing the piston rearward would be necessary to move away from the flutter boundary and arrest flutter. For these tests, however, antisymmetric flutter was not encountered over the range of tank weights and frequency ratios covered. The failure to locate this boundary may be attributed, first, to the fact that the uncoupled first antisymmetric bending to first torsional frequency ratio was 3.36, whereas the corresponding symmetric frequency ratio was 1.04; and second, to the inability of the control surfaces on the tail and the soft torsional spring inside the fuselage to keep the model positioned in a near-horizontal attitude sufficiently well to permit satisfactory tests at Mach numbers higher than 0.80 for the fuselage free-to-roll condition.

The term symmetric flutter as used in this paper applies to flutter where the wings were moving in phase and was obtained when the fuselage was locked to prevent roll. However, from observation and motion pictures made during the tests, together with oscillograms of a strain gage on the sting and an accelerometer on the model center line at the wing elastic axis, it is clear that vertical motion of the fuselage was present. This motion was particularly noticeable for more violent flutter responses in which the entire model and the sting appeared to be pitching about some axis in the sting support system. Subsequent static and dynamic studies of the sting support system revealed that the sting has considerable flexibility and that an axis of rotation exists in part of the supporting structure just beneath the tunnel floor. Some information on the elastic and inertial properties of the sting support system is given in the appendix. Furthermore, the fundamental natural frequency of the entire sting model system appears to be well within the range of the model frequencies. These dynamic characteristics of the sting mount should be kept in mind when interpreting the symmetric flutter results obtained in this investigation.

Remarks on model flutter results and related full-scale flight flutter experience. Under certain special tip-tank inertial loading conditions involving the use of lead weights in the tip tanks, flutter was experienced in flight tests conducted by the Wright Air Development Center on the fighter-type airplane referred to in the preceding sections of this paper. Some uncertainties existed in the actual inertial loading conditions, particularly with regard to the distribution of fuel within the wing. Exploring such conditions on the model would have necessitated removal of the flutter-arresting device and was, therefore, postponed until the end of the experimental program. However, an attempt to reproduce the flight flutter point with the model would have had uncertain value since the inertial properties of the model wing did not duplicate the scaled values of the airplane wing. (See figs. 4 to 6.) Furthermore,

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Calculations and Correlation of Experiment and Theory

Methods of calculations. The calculated flutter speeds and flutter frequencies presented in this paper were determined by application of a Rayleigh-Ritz type of analysis in which a combination of uncoupled vibration modes was employed to represent the flutter mode. By use of the measured wing inertial and stiffness properties given in figures 4 to 8, these modes were computed by means of the iterative method presented in reference 15. The uncoupled cantilever first bending and first torsional frequencies computed by this method are shown as functions of μ and $I_{\alpha t}$ in figures 24 and 25, respectively. The corresponding mode shapes are shown in cross-plot form in figures 29 and 30. The uncoupled cantilever second bending and second torsional mode shapes and frequencies are given in figure 31. Figure 32 shows the uncoupled antisymmetric first bending mode shapes and frequencies.

In addition to approximating the flutter mode by means of a limited number of uncoupled modes, the flutter calculations also involved the following simplifying assumptions:

(1) The oscillating aerodynamic forces and moments acting on the wing were those derived by Theodorsen in reference 13 on the basis of two-dimensional incompressible potential-flow theory (except for one case involving compressibility noted later). The effect of wing taper was accounted for in two different ways, one of which amounted to weighting or grading the aerodynamic forces and moments according to the taper ratio as recommended in reference 16. The station at three-fourths of the span was arbitrarily chosen as a reference station and the

F+iG parts of the aerodynamic coefficients for k at this reference station were held constant along the span, whereas the remaining parts of the aerodynamic coefficients were allowed to vary with span according to the semichord ratio b/b_r . This method is for convenience identified herein by the term "graded coefficients." In the other method, labeled "constant coefficients," the reference station was also at three-fourths of the span but all parts of the aerodynamic coefficients were held independent of span for k at this station.

- (2) Aerodynamic forces and moments on the tip tank were not included.
- (3) The effect of structural damping in bending or torsion was taken into consideration in the manner recommended in reference 17; that is, it was assumed that $g_{h_1} = g_{\alpha_1} = \text{Constant}$. For most of the calculations this constant was taken as zero; but for a number of cases, a value of 0.035 was used. This value was chosen on the basis of the experimental vibration data and has already been mentioned in the section on scale factors.

With these assumptions included in the analysis, flutter calculations were performed for three different wing-root boundary conditions, two of which were symmetric and the third antisymmetric to correspond, respectively, to the locked and unlocked fuselage in the flutter experiments. The first of these boundary conditions was the ideal cantilever-wing condition, whereas the second condition was that the pitching motion of the sting was approximated by means of a simple spring-inertia system whose properties are given in the appendix and was introduced into the flutter calculations as an additional elastic degree of freedom. For the antisymmetric condition, corresponding to the unlocked fuselage, the calculations included the first antisymmetric wing bending modes given in figure 32 for two different fuselage and tip-tank rolling moments of inertia. The entire program of flutter calculations has been classified, according to the degrees of freedom used, by the code designation given in the following table in order to facilitate identification of important features in specific applications of the analysis:

Identification	Boundary	Aerodynamic assumptions		Degrees of freedom
of analysis	condition	for wing taper	Fuselage	Wing
A-1		Graded coefficients		First bending, first torsion
A-2		Constant coefficients		First bending, first torsion
A-3	Symmetric (cantilever)	Constant coefficients	None	First bending, first torsion Second bending, second torsion
A-4		Constant coefficients (compressible)		First bending, first torsion
B-1	Symmetric (sting flexible)	Constant coefficients	Fuselage pitch	First bending, first torsion
C-1		Constant coefficients	Fuselage	First asymmetric bending, First torsion
C-2	Antisymmetric	Constant coefficients	free to roll	First asymmetric bending including tip-tank rolling inertia, First torsion

For convenience in comparing the calculated and experimental results, the analytical identifications given in this table are repeated in table II. In studying the effects of structural damping as stipulated by assumption (3) given previously, only analysis A-1 was used.

The calculated flutter speeds and flutter frequencies are given in table V, which is divided into six parts, each part being identified according to the code given in the preceding table. Tables V(a), V(b), and V(c) apply to the ideal cantilever-wing boundary condition. Calculations using approximations to the flexibility in the model mounting system are given in table V(d), and antisymmetric flutter solutions are given in tables V(e) and V(f).

Results of analyses A-1 and A-2.— Comparison of the results given in table V(a) for analysis A-1 (for zero damping) with those given in table V(b) for analysis A-2 for the same tip-tank inertial conditions indicates that the use of constant aerodynamic coefficients yielded flutter speeds that were as much as 30 percent higher than those obtained with graded coefficients. This trend toward higher flutter speeds due to the former and simpler application of two-dimensional aerodynamic coefficients is similar to that found for unswept, cantilever bare wings of considerably lower frequency ratios and somewhat higher taper ratios than the frequency and taper ratios of the present configuration. (See, for example, ref. 18.)

The results of analyses A-1 and A-2 without damping are compared, on the basis of Mach number, with experimental flutter results in figures 13 to 18, except for figure 17 in which the effect of damping is also shown. Although for some cases the comparison suffers for lack of a more complete fix on the experimental flutter boundaries, notably in figures 17 and 18, it is nonetheless evident that the disagreement between experiment and calculation was enlarged as the ratio of uncoupled bending frequency to torsion frequency approached unity. Furthermore, there were consistent tendencies for the calculations to be excessively conservative when the tip-tank centers of gravity were within 1 inch to the rear of the elastic axis, and for experimental and calculated flutter speeds to converge on one another for tip-tank centers of gravity further to the rear of the elastic axis. As was pointed out in reference 11, this behavior may be due to the failure to include some estimated tip-tank aerodynamic forces and moments in the calculations. (See assumption (2) in the preceding section.) Figures 12 and 16 show that flutter speeds calculated by using constant coefficients were in better agreement with experiment than were those calculated by using graded coefficients.

Effects of structural damping.- Table $\mathtt{V}(\mathtt{a})$ and figure 17 show some effect of introducing a damping coefficient of 0.035 into the simple twomode flutter calculations involving graded aerodynamic coefficients (analysis A-1). The increases in flutter speeds, due to damping, were generally greater for tip-tank centers of gravity within 1 inch to the rear of the elastic axis than for tip-tank centers of gravity further to the rear. Case 40 for $e_{t} = 0$ shows that the flutter-speed increase was very abrupt. Moreover, it can be noted in this table that the increase in flutter speed also tended to be greater as the uncoupled frequency ratio $\omega_{\text{h}_1}/\omega_{\alpha_1}$ approached and exceeded unity. (Compare, for example, case 34 with case 47 and case 35 with case 49 for approximately the same values of μ and e_{+} .) This trend is in general agreement with that found in the extensive theoretical flutter studies of reference 19. The effects of this small damping coefficient materially improved the agreement between the calculated and experimental results, particularly for cases where the calculations without damping yielded excessively conservative answers.

Effects of higher uncoupled modes and compressible aerodynamic coefficients.— Table V(c) corresponding to analysis A-3 shows that the introduction of various combinations of higher uncoupled modes into analysis A-2 produced a negligible effect on flutter speed and flutter frequency for the case considered, namely case 64. On the basis of the results of reference 5, it is reasonable to expect a greater influence of higher modes for conditions of greater mass unbalance (corresponding to high $e_{\rm t}$ values in the present investigation) than the mass unbalance of case 64.

Although, as previously noted, most of the flutter calculations presented in this paper were based on incompressible-flow theory for the aerodynamic approximations, the effects of compressibility were explored by means of analysis A-4, which was based on constant coefficients corresponding to a Mach number of 0.7 as given in reference 20. The calculations made by using this analysis were also applied to case 64, and the results, listed for case 65 of table V(b), indicate that the flutter speed was not appreciably affected here by compressibility. As an approximate check on this effect, examination of figure 1 in reference 20 showed that the κ and frequency ratio of the model for this calculation fell approximately in a region where the same flutter speeds could be obtained with either compressible (M = 0.7) or incompressible (M = 0) coefficients.

The case chosen for these studies was one for which there existed a wide disagreement between theory and experiment and it is apparent that no improvement was obtained by introducing either higher uncoupled modes or compressible aerodynamic coefficients.

Note on coupled modes of vibration. Along with the other frequencies in figures 19, 20, and 23 are shown also the first two coupled cantilever frequencies which were calculated from the flutter stability determinants with the aerodynamic terms omitted. These computations were performed for the purpose of estimating how much of the discrepancy between measured frequencies and calculated uncoupled frequencies was due to coupling in the system. As may be observed, coupling constitutes a relatively small part of this discrepancy, especially for the first mode. The remaining gap, of as much as 30 percent, between measured and calculated coupled frequencies is probably due in large measure to the fact that the measured frequencies include a strong influence of the sting support system.

Effect of sting flexibility.— The effect of flexible pitching in the sting mount, as simulated by means of analysis B-1, is shown in table V(d) and figure 33. Case 64 was used in this study which included a range of sting frequencies $\,f_{\rm S}\,$ and equivalent sting-fuselage pitching moments of inertia $\,I_{\rm fs}\,.\,$ The calculated flutter speed is seen to be sensitive to frequency $\,f_{\rm S}\,$ but relatively insensitive to inertia $\,I_{\rm fs}\,.\,$

In order to compare the results of analysis B-1 with experiment, estimates of the parameters ($I_{f_S} = 26.4$ ft-lb-sec² and $f_S = 16.8$ cps) are made in the appendix and are shown as dashed lines on figure 33. By using the estimated values of the parameters, it can be observed that there is very little effect of sting flexibility on flutter speed for the case considered. However, the large effect of frequency f_S indicates that sting flexibility may be important when the flutter frequency is close to the sting frequency.

Antisymmetric boundary condition .- The results of the antisymmetric flutter calculations are given in tables V(e) and V(f) for analyses C-1 and C-2, respectively. These calculations involved the first uncoupled antisymmetric bending-mode shapes and frequencies presented in figure 32. The flutter speeds and flutter frequencies are considerably higher than those obtained from cantilever analysis A-2 for the same tip-tank weight and moment of inertia because of the greatly increased frequency ratios based on the antisymmetric bending frequency. An overall picture of the effect of frequency ratio in relation to tip-tank center-of-gravity location is shown for three frequency ratios in figure 34 and corresponds to analysis C-1. As may be seen, the frequency-ratio effect was considerably more pronounced for forward tip-tank centers of gravity than for rearward tip-tank centers of gravity. Notice also the reversed effect of this parameter on flutter speed as the tip-tank center of gravity is moved from a position forward of the elastic axis to a position rearward of the elastic axis. The value of I_{fr} used in analysis C-1 was estimated during the design of the model. The higher value of Ifr used in analysis C-2 was based on measured data obtained after completion of the model. However, the differences in flutter speeds between tables V(e) and V(f)are attributed more to the effect of I_{tr} than of I_{fr} . This contention is based on the existence of a small difference between the antisymmetric bending frequencies due to the two different values of I_{fr} , as compared with the larger difference caused by modifying the tip bending-moment boundary condition from zero to a finite value determined from the estimated tip-tank rolling inertia about the wing tip. (See fig. 32.) However, as comparison between tables V(e) and V(f) shows, the effect of this parameter on flutter speed was small.

The theoretical antisymmetric flutter trend plotted in figure 16 (for analysis C-2) serves to show the location of the experimental symmetric flutter region. As is implied in figure 34, this region might well have fallen within the speed range covered by the experiments, particularly for forward tip-tank centers of gravity, had the ratio of uncoupled antisymmetric bending frequency to torsion frequency been closer to unity. Preliminary analytical work during the construction of the model indicated that this condition might be realized for tip-tank weights approaching an empty condition, a condition which was not attainable in this investigation.

CONCLUSIONS

An extensive investigation has been made of flutter of a true-speed dynamically scaled model representative of unswept-wing fighter-type airplanes with tip tanks. The model was one-sixth the size of an unswept

fighter airplane having an aspect ratio of 6 and was of spar-balsa segment-type construction with a flutter-arresting device in each tip tank. The following conclusions can be drawn from the results of this investigation:

- 1. A true-speed dynamically scaled model can be built to satisfy a given set of elastic and inertial specifications for flutter tests at high subsonic speeds.
- 2. In general, symmetric flutter speeds were increased for tip-tank center-of-gravity positions forward of the elastic axis and tended to decrease as the ratio of uncoupled wing bending frequency to torsion frequency f_{h_1}/f_{α_1} approached unity.
- 3. Increasing the volume of the tip tanks by a factor of 3 tended to decrease the flutter speed. The installation of a horizontal fin at the rearward end of the tank may have had a beneficial effect on the flutter speed for the small tank but may have had a slightly detrimental effect for the larger tanks.
- 4. Symmetric flexure-torsion flutter was effectively arrested when the tip-tank center of gravity was quickly shifted forward of the elastic axis. For cases where the center of gravity could not be shifted this far forward, the effectiveness was not as positive and was in some cases negative. In applications to flight flutter use, the weight, size, and time factors involved in the effectiveness of this device must be considered.
- 5. Calculated flutter speeds based on assumptions involving structural damping, higher structural modes, compressible aerodynamic coefficients, and sting-support flexibility (and not including aerodynamic forces and moments on the tip tanks) were, in general, excessively conservative for tip-tank centers of gravity near the elastic axis. However, the introduction of structural damping improved the agreement between calculation and experiment and resulted in greater increases in calculated flutter speeds, especially for ratios of uncoupled wing bending frequency to torsion frequency near unity. The effects of sting-support flexibility showed that this modification to the calculation may be important where the natural frequency of the sting system was near the flutter

frequency, and further work on these effects is in order. In addition, this sting flexibility should be kept in mind in interpreting the experimental flutter results.

Langley Research Center,
National Aeronautics and Space Administration,
Langley Field, Va., September 29, 1954.

APPENDIX

CHARACTERISTICS OF MODEL MOUNTING SYSTEM

Description

The model was mounted in the Langley 16-foot transonic tunnel in the manner shown in figure 1(a). This mounting system essentially consisted of two main parts, the sting and the sting-support structure. The juncture between these two parts was 95 inches rearward of the elastic axis along the tunnel center line, and the actual connection was made by means of bolts approximately 1 inch in diameter located around the periphery. The sting had a joint whose center line was located 60 inches rearward of the elastic axis along the tunnel center line. At the free or upstream end of the sting a sting extension was mounted which was part of the main structural member of the model fuselage.

The sting support structure was made up of a massive streamlined strut mounted to a circular-shaped track that was part of the mechanism used to provide a change in angle of attack on models in the tunnel. The top of the strut to which the base of the sting was connected was cylindrical in shape with the longitudinal axis of the cylinder lying along the tunnel center line.

Stiffness Distribution

Figure 35 gives the stiffness distribution along the length of the sting. The variation shown is based on both calculation and experiment, the experimental values being obtained from measurements of the bending slope with the sting subjected to vertical loads applied to the main structural member of the fuselage at the elastic axis. Along the solid portion of the curve the calculated and experimental values were in close agreement. The dotted portions of the curve apply to the various joints and connections in the system and represent the most reasonable stiffness values corresponding to the actual measured slopes over these discontinui-The experiments further revealed that the sting support structure was, for all practical purposes, rigid in the vertical or pitch direction, but that there was a finite stiffness contribution from the strut mount on the circular track. An estimation of this stiffness in terms of a spring constant for the pitching degree of freedom is 596×10^6 in-lb/radian, based on the approximate location of the actual axis of rotation in the base of the sting support structure.

Weight Distribution

Figure 36 shows the variation of weight per unit length with the distance along the tunnel center line. This distribution was calculated from the specifications of the sting.

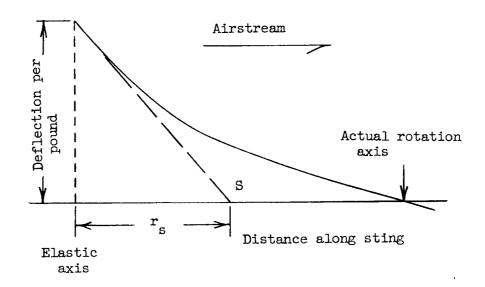
Fuselage Inertial Properties

The weight and pitching moment of inertia of the fuselage, including the main structural center section, wing spar root brackets (see fig. 3(a)), nose, center, and rearward fairings with control-surface assembly, are given as follows:

Fuselage weight, lb	157
Mass pitching moment of inertia of fuselage about	
elastic axis, in-lb-sec ²	75.3
Fuselage center of gravity, rearward of elastic	
axis, in	5.92

Simplified Representation of Sting Flexibility

A scheme is advanced herewith for approximating the actual elastic behavior of the sting mount by means of an equivalent one-degree-of-freedom system oscillating in pitch about a point S in the manner shown in the following sketch:



The equivalent elastic properties consisted of a spring constant $k_{\rm S}$ and an effective rotation arm $r_{\rm S}$. These parameters were determined from static loadings at the elastic axis together with deflection-curve measurements along the sting. As indicated in the foregoing sketch, the slope of the deflection curve at the elastic axis established the point S as being the effective axis of rotation and $r_{\rm S}$ as the effective rotation arm. From these measurements the following values were found for these two parameters:

$$k_s = 4015 lb/in$$
.

and

$$r_s = 42 in.$$

In addition,

$$I_{f_S} = 26.4 \text{ ft-lb-sec}^2$$

which is based on the measured fuselage inertial properties given previously and includes the effective sting mass corresponding to $r_{\rm S}$ = 42 inches.

On the basis of the foregoing structural and inertial properties, the estimated frequency of the fuselage-equivalent sting combination is

$$f_s = \frac{1}{2\pi} \sqrt{\frac{k_s r_s^2}{2^4 I_{f_s}}} = 16.8 \text{ cps}$$

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TABLE I .- GEOMETRICAL AND INERTIAL PROPERTIES OF MODEL TIP TANKS

(a) Small tip tank: tank A (scaled 230-gallon tip tank)

[Length, 30.3 in.; maximum diameter, 4.24 in.; span of fin (from tank center line), 4.38 in.; exposed area of fin, 11.2 in.²]

Piston		Without	fin		With	fin
weight,	W _t ,	e _t , in.	$I_{lpha_{ extsf{t}}},$ in-lb-sec ²	W _t , lb	e _t , in.	Ι _{ατ} , in-lb-sec ²
0.40	2.70	1.46	0.280	2.78	1.96	0.332
.78	3.08	2.49	.378	3.16	2.92	.430
1.07	3.37	3.22	.466	3.45	3. 59	.520

(b) Large tip tank: tank B

[Length, 43.7 in.; maximum diameter, 6.12 in.; span of fin (from tank center_line), 6.31 in.; exposed area of fin, 23.9 in.²]

Piston		Without	fin		With	fin
weight,	W _t ,	e _t ,	Ι _{ατ} , in-lb-sec ²	W _t , lb	e _t , in.	$I_{\alpha t}$, in-lb-sec ²
0.40	4.77	1.22	0.698	5.00	1.92	0.803
.78	5.15	1.73	.680	5 .3 8	2.48	.901
1.07	5.44	2.22	.769	5.67	2.92	.990

Table	Description
III	Experimental - comparable natural frequencies of model and a full-scale fighter configuration
IV V	Experimental - results of model flutter program Analytical - results of flutter calculations

Figure	Tank size μ		f_{h_1}/f_{α_1}	e _t , in.	Wing-root boundary condition	Analysis
			Ordinate,	M; abscissa,	et	
12	A and B	0.60	0.92	-4.9 to 1.7		None
13		•59	.84, .93	-4.9 to 3.5		A-1, A-2
14	A	.82	.90, .96, 1.02	-3.1 to 4.0	Fuselage clamped	A-1
15		1.005	.92	-2.4 to 4.0	_	A-l
	·		1.04	-2.4 to 4.0		A-1, A-2
16		1.02	3.36	-2.6 to 1.6		None
	В		3.36	0 to 2.4	free to roll	C-2
17		.895	. 98	-3.8 to 4.0	Fuselage	A-1
18		1.29	1.01, 1.09	-2.1 to 4.0	clamped	A-1
			Ordinate,	f; abscissa,	et	
			.84	-1.0 to 2.8		A-2
19		•595	.84	-2.0 to 3.5		First two coupled frequencies
			-93	-4.8 to 1.7		A-2
20	A	•59	-93	-1.25 to 1.6	Fuselage clamped	First two coupled frequencies
21		.82	1.02	9 to 2.5	Clamped	A-1
22		1.005	.92	-2.3 to 4.0		A-l
			1.04	-2.4 to 4.0		A-1
23	В	1.02	1.04	0 to 3.2		First two coupled frequencies

TABLE III. - COMPARISON OF MODEL VIBRATION DATA WITH SCALED VIBRATION DATA OF A

COMPARABLE FIGHTER-TYPE AIRPLANE WITH TIP TANKS

19.6 29.5 18.9 31.2	.506 .401
	5 ⁴ 4.
る 年 45万	0

I Model mounted on dummy sting for vibration survey. (See fig. 1(b).)

Model mounted on sting in tunnel. (See fig. 1(a).) Data from test 5 in table IV. H

Airplane in level attitude resting on its main landing gear with tires at normal pressure and fuel in integral wing tanks. 4

Airplane, with empennage and afterburner removed, inclined at 190 angle of attack with landing gear on chocks and blocks, tires at normal pressure, and fuel in integral wing tanks. ф

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TABLE IV .- RESULTS OF EXPERIMENTAL FLUTTER PROGRAM

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- Phase	rinos angle, Remarks deg	i	No flutter within test range; bursts of torsional oscillations from approximately $M=0.60$ to $M=0.80$	No fluther within range of available tunnel power	No definite flutter point at available tunnel pover; flutter in bursts over vide range of Mach number below top point	7 0 Lociplent fluther	No flutter within range of available tunnel power	Incipient flutter; piston relatively ineffective	Very close to flutter; formard center-of-gravity shift due to firing piston appeared to be insf- fective in arresting flutter	6 5 Symmetric flutter involving fuselage translation	No flutter	3 48 Plutter	5 % Plutter	No flutter; near to flutter at M = 0.73	Repeat of test 12; burst of flutter at M = 0.73; test continued to M = 0.81 without obtaining flutter	5 10 Plutter	3 37 Flutter	No flutter	No flutter; random oscillations appeared to be slightly worse with piston forward	55 Flutter	Test discontinuad; right tip tank loose	by Flutter
Ī		gd _D	·			21.7		8/8 7:4:	8.8	17.8		19.8	39.3		18.6	16.5	15.8			15.0		16.0
	cps	After flutter	3. 04.		57.0	36.9× 4.6.	87. 27.85		ないま	28.7						18.2						
Frequency data	ftl, cps	Before After flutter flutter	41.14	38.5	38.54 4.0.54	4.4. 28.	28.5	¥.5	88	23.5 23.8 29.8 29.8	21.7 26.65	21.3	22.3 18.95	88. 7.	26.6	18.2	8.50 19.0	18.5 19.95	18.5 19.95	19.05 18.55	19.45	19.35
Preq	sď.	Before After flutter flutter	19.25		17.4	17.4 17.65	17.65 18.1		15.4	15.2						14.6 14.2						
	f _{bl} , cps	fore																				
		AG	20.1	17.95	16.1 16.7	18.5 18.45	18.2 17.8	15.0	7.4.1 7.4.1	14.6	16.4	15.8 13.95	16.1	15.3	15.1	14.6	1.41	25.55 7.51	2:42 12:52	13.6 14.6	5.4. 5.4.	4.51
	æ	er foot							.,													_
mamic data	6,4	ft per foot	4.2 × 106	.001674 4.0 17.85	1.61 4.0 16.1	.001.706 3.8 18.5	.001546 4.0 18.2	.001852 3.5 15.0	7.41 3.6 at8100.	.001859 3.8 14.6	.001388 3.8 16.4	3.6 15.8 15.8 15.95 15.95	.001659 3.5 16.1	3.6 15.3	3.8	.002158 2.9 14.6	.001777 5.4	.001610 4.2 14.5	.001672 4.1 14.5	3.51 7.5 o19100.	.002096 5.2 13.5	7.51 5.5 19.4 14.7
Aerodynamic data	6,4	ft per foot	0.001454 4.2 × 106					3.5	3.6													_
Aerodynamic data	6,44	M Vr, ips 10-8ec	4.2 × 106	0.4 4.00	0.4 175100.	9.6 307.100.	0.4 34₹:00.	.001852 3.5	9.5 018100.	8.5 99100.	.001388 5.8	3.6	.001859 3.5	3.6	9.5 659100.	.002158	4.6 777100.	2.4 016100.	.001672 4.1	7.5 019100.	.002096 3.2	.001991 3.5
	6,44	M Vr, ips 10-8ec	1,015 0.001454 4.2 × 106	852 .001674 4.0	0.4 175100. 658	803 .001706 3.8	0.4 942:00.	686 .001852 3.5	9.5 018100. 617	740 .001859 5.8	999 .001388 5.8	477 co1718 3.6	683 .001859 3.5	9.6 775100. 938	816 .001659 3.8	164 .002158 2.9	4.5 777100. 617	2.4 019100. 016	967 .001672 4.1	7.5 019100. 599	5.2 302096 3.2	629 .001991 5.5
et, in. Aerodynamic data	C 40 41 22 3	M Vr, ips 10-8ec	0.94 1,013 0.001454 4.2 × 106	.78 852 .001674 4.0	0.4 175100. 658 87.	8.5 307.100. 608 67.	.84 922 00. 546 4.0	.62 686 .001852 3.5	9.6 018100. 617 69.	.68 740 .001859 3.8	.88 999 .001388 5.8	9.6 8tytoo. 477 69.	.61 685 .001859 3.5	3.6 772100. 338 77.	9.5 959100. 918 57.	9.2 83 ماد.	4.5 777100. 917 49.	5.4 o100. o16 48.	.80 867 .001672 4.1	7.6 019100. 599 49.	.50 50200. HH2 03.	.57 629 .001991 3.5
Ĥ.	Platon w 623 At 2002	forward N.f., 198 10-8ec	1.21 0.94 1,013 0.001454 4.2 × 106	1.92 .78 872 .001674 4.0	2.52 .78 853 .001571 4.0	.73 803 .001706 3.8	0.4 346:00. 126 48.	.62 686 .001852 3.5	2.04 .001810 .001810 3.6	81.18 .001859 3.8	86. 28. 39. 4.4.	.5.42 .001718 3.6	.2.07 683 .001859 3.5	3.5 TT21.00. 338 TT. 3.6	9.5 626100. 618 67.	-2.72 .45 491 .002158 2.9	1.08 4.17 3.4	2.4 010100. 016 48. TT.S-	-2.64 .001672 4.1	7.5 019100. 699 49. 77.1-	.50 50±44 .002096 3.2	62 .001991 3.5

Petter designation denotes tip tank used; the set of numbers following first dash refers to µ; the set of numbers following second dash gives the uncoupled frequency ratio fh_1/fo_1.

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TABLE IV .- RESULTS OF EXPERIMENTAL FLUTTER PROGRAM - Continued

	Remarks	Violent flutter; symmetric mode involving fuselage translation	No flutter	Plutter	Plutter	Visual evidence of flutter; piston in forward position appeared wery effective in arresting flutter	No flutter; torsional oscillations appeared stronger and less damped with piston forward than with piston rearward	Wery close to flutter; piston appeared ineffective in arresting flutter	No flutter within range of test; fitting for left wing tank loosened during test	Pluter	Flutter	Plutter	Plutter	Plutter	Flutter	Plutter	Flutter in bursts	No flutter within range of available tunnel power; action of piaton wery effective in preventing flutter	Flutter, action of platon very effective in preventing flutter	Flutter; piston effective in stopping flutter	Plutter
5	angle, deg	27		ĸ	15	Ş					*8	Ж	27	8	ex.		ส		87	88	31
	r cps	15.4		15.6	<u>2</u>	16.7				15.1	<u> </u>	13.5	21.5	19.4	8.	21.0	16.6		15.6	15.0	14.5
late	l, cps After er flutter	19.0		17.4		18.5	17.7			16.5	16.95			22.2			17.9			17.4	17.2
Frequency data	ftl, Before r flutter	25 g			18.05	18.0	18.9	16.5		16.9	17.2	18.4	88	28 7.6.	21.5	23.1	17.65	17.6 19.9	19.7	16.9	17.45
Pre	f _{b1} , cps Before After flutter flutter	1.54		13.9		13.0	25.0 8.0		ы.6	12.3	12.65		14.95	14.7			13.1			15.75	25.51 2.52
		12.4	_		15.1	6.51 12.9	일 점 6.69	†∙टा	8.डा	13.3	11.8	8.55 6.51	14.95 14.8	14.7	7.41	7.41	13.05	13.4	13.0	13.0	12.9 13.2
1c data	lb-sec2 R,	0.001951 5.7 × 106	.601516	.001875 5.7	.002132 3.1	0.4 797100.	.001683 4.2	.002012	.001587 3.9	.001871 3.6	4.5 646100.	.001896 3.5	.001882 3.7	.001985 3.4	4.5 156100.	.001955 3.3	.002082 2.6	.001701	5.6	.002124 2.4	.002211 2.7
Aerodynamic data	al s _d														8	8	90.	8	8	00.	<u>8</u>
Ye.	, G	0.63 689	.85 939	601 59.	716 74.	.73 790	.81 872	.57 574	.81 890	.63 692	7 632	1 668	2 708	9	.58 645	929 9	1 60	<u>&</u>	9 759	8 428	0 4.37
	ston. Rerd M	0.t.7			- 7 .						15.	.52 51	- 64	 		δ χ	-t ₁	-18 -18	- -	& &	₹.
e _t , in.	Piston Piston rearward forward	, 8 8	38	1.05 -1.64	1.83	.15 -2.45	.32 -2.28	.72	10 -2.31	97.	0219		13 -1.49	99 -1.52	.02	.9. o4	18 -1.43	84 84.	%: 44	07.	-99
	Pis	ai 	·	-i 	-i		•		<i>i</i>	••	2.02	2.53	13	1.09	· —	<u>~</u> .	1.18	-1. X			
	Configuration (a)	A82597	A82-1.02	A82-1.02	A82-1.02	A8595, no fin	A8695	A90-1.01, no fin	A-1.00592	A-1.00592	A-1.00592	A-1.00592	B59594	B60593	в62590	B62590	B8598; no fin	в89598	в89598	B85-1.02, no fin	B875-1.04, no fin
	Test	12	83	23	77.	ĸ	92	27	28												

*Letter designation denotes tip tank used; the set of numbers following first dash refers to u; the set of numbers following second dash gives the uncoupled frequency ratio full/fag.

TABLE IV. - RESULIS OF EXPERIMENTAL FLUTIER PROGRAM - Concluded

	Renarks		Flutter	No flutter; wing showed strong bending tendencies	Flutter; strong torsion response	Flutter	Flutter; amplitude of oscillations increased slowly prior to flutter	Flutter	No flutter; response of model observed during test appeared as close to flutter with piston forward as with piston rearward	Test discontinued; no flutter observed	No flutter; predominant bending response	Sustained flutter; predominantly torsion; piston very effective	No flutter; model free to roll	No flutter; model free to roll; very unstable in roll; records indicate model less damped with piston in forward position near top Mach numbers	Flutter; model free to roll; piston forward effective in stopping flutter	Flutter in bursts; model free to roll; strong torsion response
	Phase angle	deg	87		8	43		£				£4			53	
	-	r cps	14.8		14.3	14.2	14.0	13.3				7:21			14.6	14.3
	ftl, cps	After flutter	16.85 16.7			15.8		4.51				13.4 13.3				
ata	ft],	Before flutter	16.9 16.85	15.0	15.2	14.1 15.55	15.3	15.55	16.5	13.5	8. य	13.4	14.8	15.1	14.8	4.41 9.41
Frequency data		After flutter	7.21 2.61			5. 5. 5. 5.		н.3 г.5				1.10				_
Fre	f _{b1} , cps	Bcfore flutter	12.4 12.6	11.7 11.0	11.3 11.6	12.1 12.3	11.ò 12.15	11.3 11.3	10.7 11.14	10.65 10.5	10.1	8.01 4.11	13.3 (symmetric) 43.7 (antisymmetric) 16.4 (symmetric) 44.6 (antisymmetric)	13.2 (symmetric) 42.3 (entisymmetric) 43.0 (antisymmetric)	13.1 (symmetric) 42.4 (antisymmetric) 12.9 (symmetric) 42.4 (antisymmetric)	12.5 (symmetric) 42.2 (antisymmetric) 12.9 (symmetric) 43.4 (antisymmetric)
	α	per foot	2.4 × 106							-7.7						
Aerodynamic data	p, 2		0.002255 2.	9.5 446,100.	.001372 3.5	.002213 2.6	.002232 2.5	.002135 2.7	-		.001766 3.7	.002223 2.6	.0017100	.001900 3.5	.001912	.002179 2.6
Aerodyna	-	'f', fps	396	686). TS#	, t14	7. 2.14		•		, e e e e e	.67	701	0. 699	0. 0444
	<u> </u>	M V _f	0.36	<u>a</u> ġ.	- 68-	<u> </u>	.37	24.	.37	۲۲.	- 169:	4 65.	7.	8.	8	91.
	ptetos	orward	-1.33	-2.46	-1.54 	-1.16	-1.17		04.1	-2.08	-1.92	. 97	-2.59	-2.46	-1.64	80.
et, in.	T at a	<u></u>	1.19	1.27	. 65	1.04	1.03	1.62	3.58	57	- 25:-	92.		75	35.	1.63
	Test Configuration		B875-1.07	B-1.02-1.04	B-1.02-1.04	B-1.01-1.04	B-1.01-1.05	3-1.025-1.04	B-1.025-1.04	B-1.30-1.01	B-1.31-1.07	B-1.28-1.09	B-1.025-1.04	B-1.02-1.04	B-1.02-1.04	B-1.025-1.04
	Test		-	24	£4	‡	₹ 	- - 9π		£8	64	R	51 B	SZ.	53	± ±

^aLetter designation denotes tip tank used; the set of numbers following first dash refers to μ ; the set of numbers following second dash gives the uncoupled frequency ratio f_{h_1}/f_{c_1} .

TABLE V.- RESULTS OF FLUTTER CALCULATIONS

(a) Cantilever-wing boundary condition, graded coefficients, first bending and first torsion (analysis A-1)

		g _h	= g _{αl} =	0	$g_{\mathrm{h_1}}$	$= g_{\alpha_1} = 0$. 0 3 5			
Case	e _t , in.	v_f, fps	f _f , cps	l/k _r	v _f , fps	f _f , cps	$1/k_{f r}$			
	A56591 configuration									
ı	3. 925	699	20.8	13.09						
			A5993	configura	ation					
2 3 4 5 6 7	-1.16 0 .24 .50 1.00 1.59	501 443 461 519 558	22.7 22.6 22.5 22.2 21.8	8.60 7.66 8.00 9.12 9.96	562.5 523 547 583.5	22.25 22.2 22.1 21.95	9.85 9.18 9.63 10.35			
	A5984 configuration									
8 9 10 11 12 13	0 .50 1.00 1.75 2.64 3.50	993 774 748 761 785 803.5	24.1 23.7 23.3 22.6 21.8 21.1	16.05 12.74 12.51 13.13 14.04 14.84	849 774 783 831 886	23.2 23.0 22.7 22.3 21.9	14.28 13.1 13.43 14.54 15.75			
		Į.	A815 - .90	configura	ation					
14 15 16 17 18 19	0 1.00 2.00 3.00 3.695 4.00	655 561 615 652 667 672	19.9 19.4 18.6 17.9 17.4 17.2	12.82 11.26 12.9 14.2 14.95 15.23	582 658.5 729 777 781	19.3 19.0 18.5 18.3 18.0	11.7 ⁴ 13.51 15.31 16.56 16.9			
		Į.	182596	configura	ntion					
20 21 22 23 24 25	0 .50 1.02 1.97 3.00 4.00	329 367 442 516 563 592	18.9 18.8 18.5 17.9 17.2 16.6	6.78 7.60 9.32 11.23 12.75 13.9	422 475•5 574 656	18.65 18.6 18.3 17.9	8.80 9.98 12.23 14.27			

TABLE V.- RESULTS OF FLUTTER CALCULATIONS - Continued

(a) Cantilever-wing boundary condition, graded coefficients, first bending and first torsion (analysis A-1) - Concluded

		e	$s_{h_1} = s_{\alpha_1} =$	0	g_{h_1}	$= g_{\alpha_1} = 0$.035				
Case	e _t , in.	v _f , fps	f _f ,	l/k _r	v _f ,	f _f ,	1/k _r				
	A82-1.02 configuration										
26 27 28 29 30 31 32	0 .50 .96 1.74 2.50 3.00 4.00	255 317 384 439 484.5 507 540	18.1 18.0 17.8 17.4 17.0 16.75 15.2	5.49 6.85 8.40 9.82 11.11 11.8 12.98	405.5 454	18.0 18.05	8.72 9.80				
			A-1.005	92 configura	ation		•				
33 34 35 36 37 38	.25 .71 1.95 2.47 3. 00 4.00	477 465 538 570 595 620	17.5 17.3 16.6 16.3 16.0 15.4	10.6 10.48 12.62 13.63 14.43 15.68	498 586 636 674	17.3 17.0 16.75 16.6	11.21 13.45 14.8 15.83				
	B89598 configuration										
39 40 41 42 43 44	-1.015 0 .69 1.50 2.50 4.00	217 • 5 376 472 536 585	18.0 17.8 17.3 16.7 15.8	4.705 8.22 10.64 12.5 14.43	655 425 531 636	17.9 17.8 17.7 17.2	14.29 9.29 11.7 14.4				
	· · · · · · · · · · · · · · · · · · ·		B-1.02-1.0	04 configure	tion	·	.				
45 46 47 48 49 50 51	27 .0135 .65 .94 1.96 3.58 3.99	537 328 362 426 483.5 497	16.3 15.9 16.0 15.6 14.5 14.6	12.85 8.03 8.81 10.65 12.96 13.28	435 525	16.15 15.95	10.49 12.84				
			B-1.29-1.0)l configura	tion						
52 53 54 55 56	35 .50 1.50 2.50 4.00	288 407 465 508	14.75 14.3 13.8 13.1	7.62 11.1 13.14 15.15	376 592	14.85	9.86 16.14				
			B-1.29-1.0	9 configura	tion						
57 58 59 60	.99 1.50 2.50 4.00	338 360 404 446	13.8 13.6 13.2 12.6	9.54 10.33 11.92 13.77	501	14.15	13.8				

TABLE V.- RESULTS OF FLUTTER CALCULATIONS - Continued

(b) Cantilever-wing boundary condition, constant coefficients, first bending and first torsion (analysis A-2)

									
Case	e _t , in.	v _f , fps	f _f , cps	1/k $_{f r}$					
	A5	6591 conf	iguration						
61	3.925	762	20.9	14.25					
A5993 configuration									
62 63 64 65a 66 67 68	-1.16 0 .24 .24 .50 1.00 1.59	664 486 488 491 524 604	22.7 22.6 22.2 22.6 22.8 21.9	11.39 8.39 8.67 8.50 8.97 10.8					
	A5984 configuration								
69 70 71 72 73 74 75	-2.00 0 .50 1.00 1.75 2.64 3.50	999 8 31 81 3 819 855 890	24.1 23.7 23.3 22.6 21.8 21.1	16.16 13.66 13.61 14.14 15.3 16.45					
	B-1.	02-1.04 cont	figuration						
76 77 78 79 80 81 82 83 84	27 .0135 .50 .94 1.50 1.96 2.75 3.22 3.99	720 374 407 442 475 511 532 556	16.2 16.1 15.9 15.6 15.2 15.0 14.7	17.35 9.05 9.95 10.86 11.86 13.08 13.82 14.71					

aComputed using Mach number 0.7 coefficients (analysis A-4)

TABLE V.- RESULTS OF FLUTTER CALCULATIONS - Continued

(c) Cantilever wing boundary condition, constant coefficients, first bending, first torsion, second bending, and second torsion

[Analysis A-3; configuration A-.59-.93; $e_t = 0.24$ in.]

Case	Identification of modes used	v _f , fps	f _f , cps	1/k _r
85	First bending, first torsion, second bending, second torsion	479	22.7	8.23
86	First bending, first torsion, second bending	476	22.6	8.20
87	First bending, first torsion, second torsion	476	22.6	8.20
64	First bending, first torsion	486	22.6	8.39

(d) Symmetric (sting-fuselage pitching) boundary condition, constant coefficients, flexible sting pitching, first cantilever wing bending, and first cantilever wing torsion

[Analysis B-1; configuration A-.59-.93; $e_t = 0.24$ in.]

		T								
Case	I _{fs} , ft-lb÷sec ²	γ _f , fps	f _f , cps	1/k $_{f r}$						
		f _s = 14.5 cps								
88 89 90 91 92	25 39.8 50 70 100	465 431 427 435 447	23.0 23.0 23.3 22.75 22.8	7 .87 7 .29 7 .13 7 .45 7 .63						
	f _s = 16.7 cps									
93 94 95 96 97 98	10 25 40 50 70 100	518 509 463 428 447 426	22.7 23.05 23.0 22.9 22.8 22.9	8.90 8.61 7.82 7.28 7.64 7.26						
		f _s = 20 cps								
99 100 101 102 103	10 25 39.8 50 70	461 658 536 528 469	22.9 23.1 23.1 23.1 22.9	7.84 11.075 9.05 8.91 7.97						
		$f_s = 23 \text{ cps}$								
104 105 106 107 108	10 25 40 50 70	995 1,112 970 990 999	23.0 23.3 24.35 23.1 23.0	16.85 18.58 15.52 16.72 16.90						

TABLE V. - RESULTS OF FLUTTER CALCULATIONS - Continued

(e) Antisymmetric (fuselage free to roll) boundary condition, constant coefficients, first asymmetric bending, and first cantilever wing torsion

Analysis C-1; configuration B-1.02-1.04;
$$I_{fr} = 0.0605 \text{ ft-lb-sec}^2$$
; $I_{tr} = 0$

Case	e _t , in.	v _f , fps	f _f , cps	l/k _r
109	-2.50	3,441	31.6	42.5
110	-2.00	4,213	31.9	51.55
111	-1.00	7,125	31.3	88.8
112	75	9,995	31.0	125.7
113	50	2,829	15.4	71.7
114	0	1,427	15.1	36.7
115	1.19	1,088	14.6	29.2
116	2.375	1,027	14.0	28.6

TABLE V .- RESULTS OF FLUTTER CALCULATIONS - Concluded

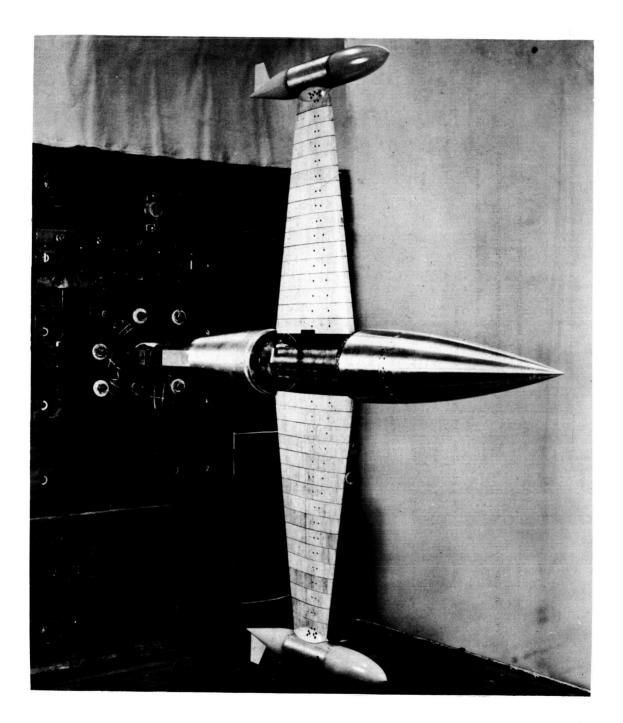
(f) Antisymmetric (fuselage free to roll) boundary condition, constant coefficients, first asymmetric bending, and first cantilever wing torsion

Analysis C-2; configuration B-1.02-1.04;
$$I_{fr} = 0.0984$$
 ft-lb-sec²; $I_{tr} = 0.0192$ ft-lb-sec²

Case	e _t , in.	v _f , fps	f _f , cps	1/k _r
117	0	1,489	15.05	38.6
118	1.19	1,131	14.4	30.6
119	2.375	1,059	13.8	29.9



L-77976.1 (a) Model in the Langley 16-foot transonic tunnel. Figure 1.- Photographs of dynamic model.



(b) Model mounted on dummy sting for vibration survey. L-77439

Figure 1.- Concluded.

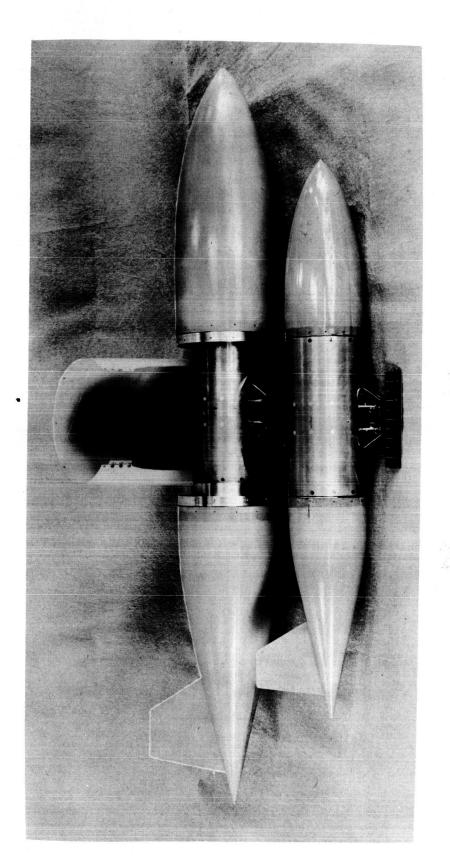
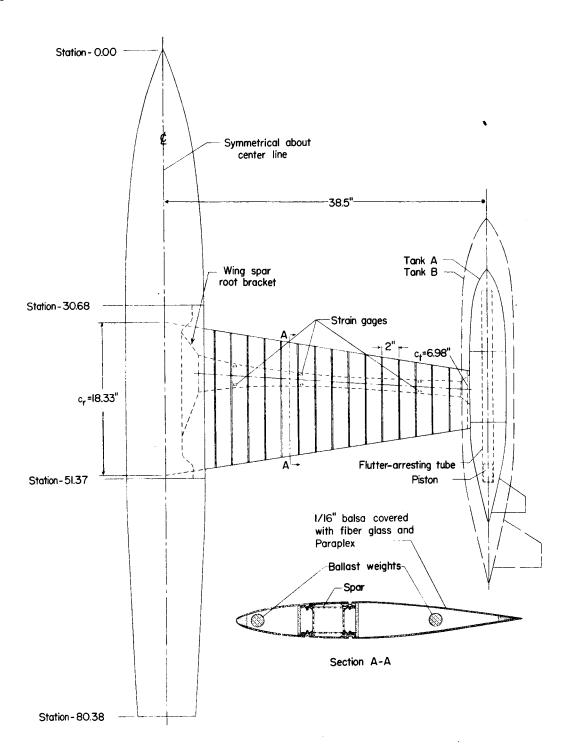
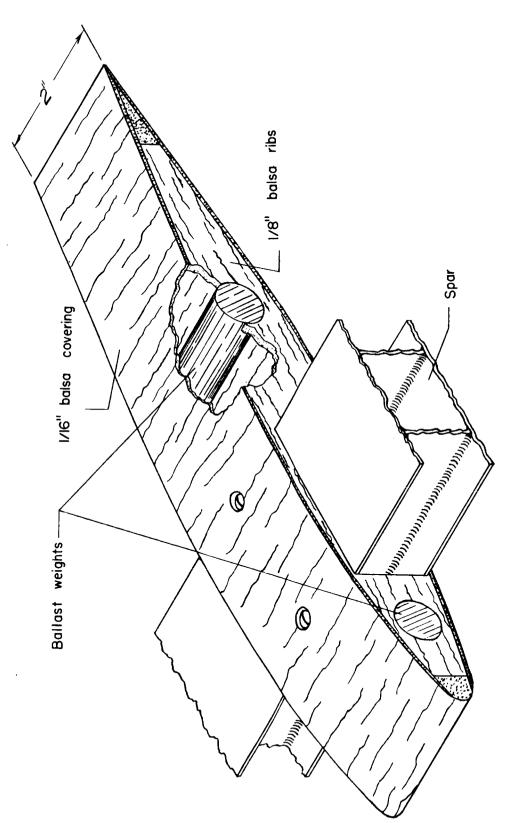


Figure 2.- Tip tanks used in investigation.



(a) Plan view of model.

Figure 3.- Sketches of dynamic model.



(b) Sketch showing balsa segment construction for a single segment.

Figure 5.- Concluded.

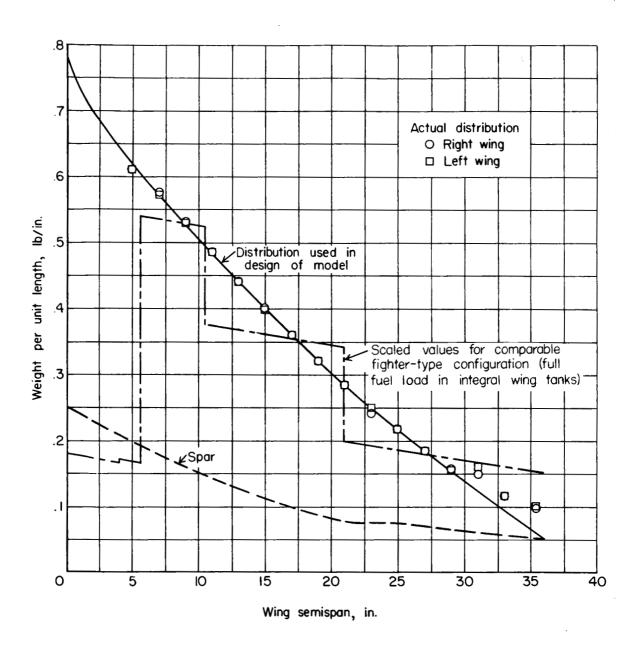


Figure 4.- Spanwise weight distribution.

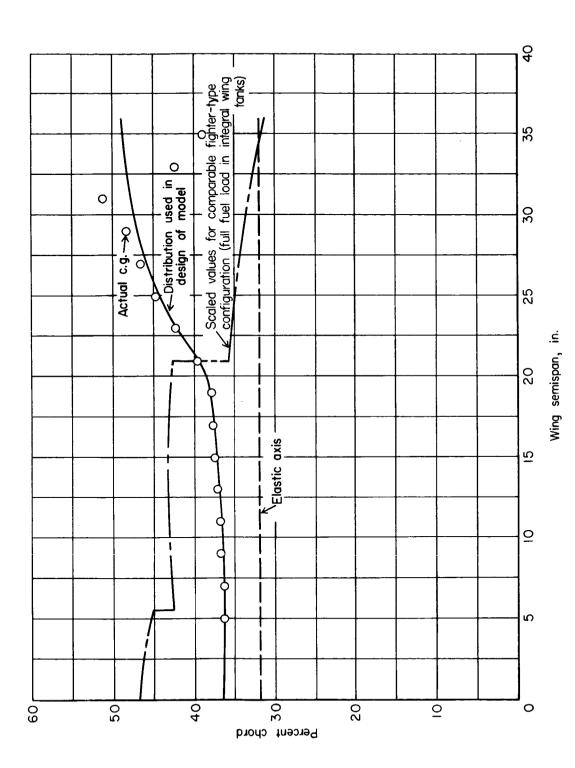


Figure 5.- Variation of center of gravity and elastic axis with span.

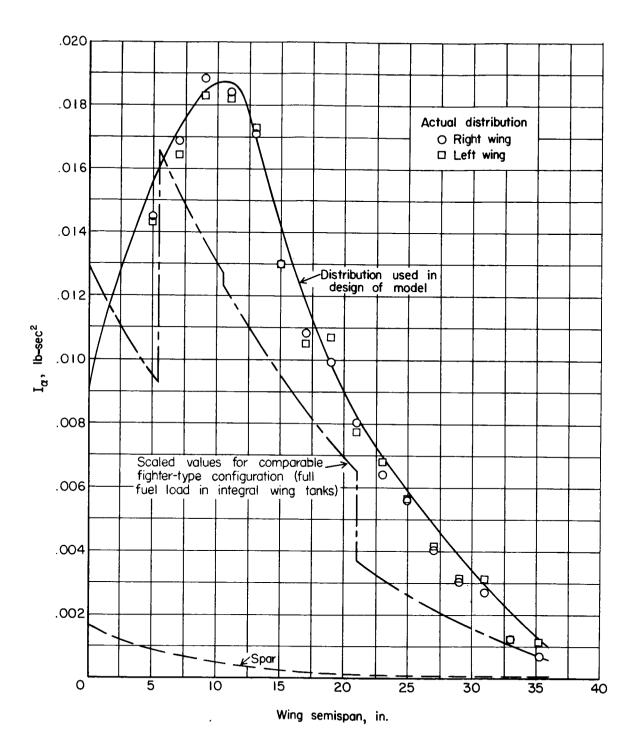


Figure 6.- Variation of moment of inertia with span.

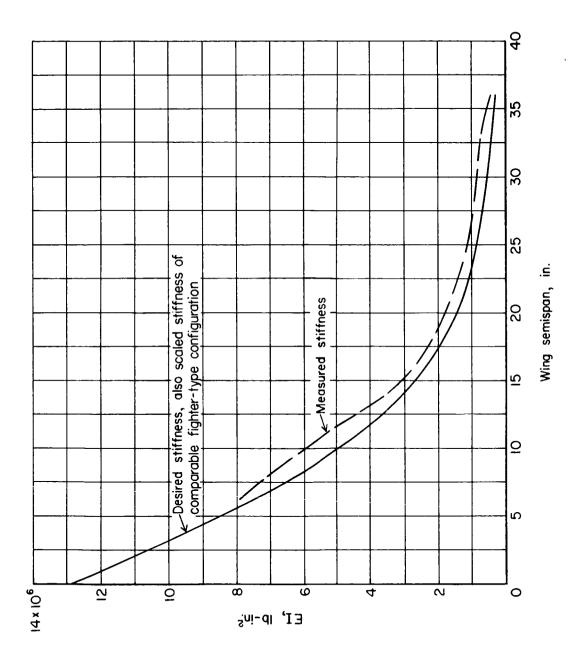


Figure 7.- Variation of bending stiffness with span.

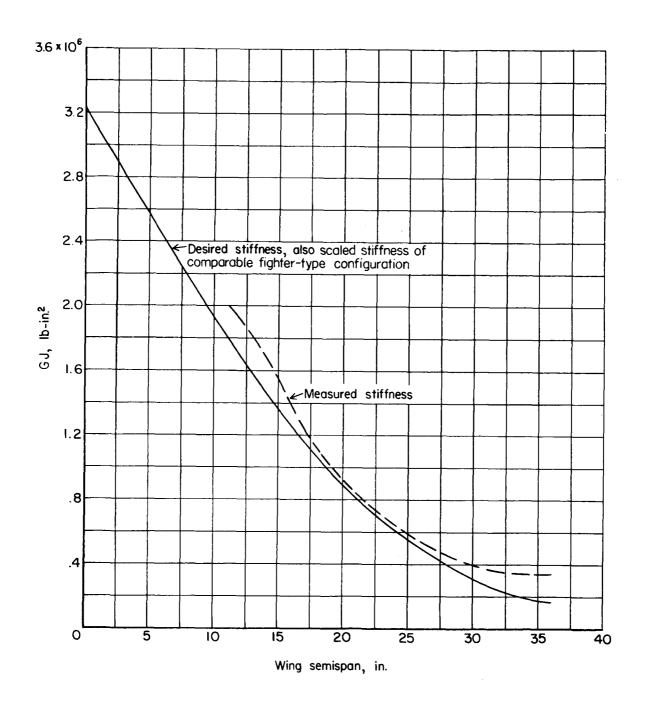
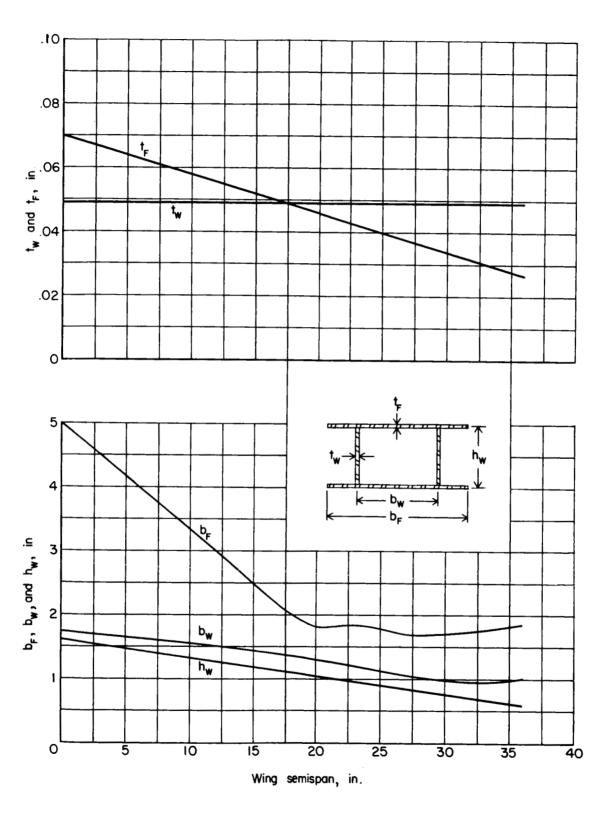


Figure 8.- Variation of torsional stiffness with span.



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Figure 9.- Design specifications for wing spar.

Figure 10.- Sketch of tip tank A showing ballast weights.

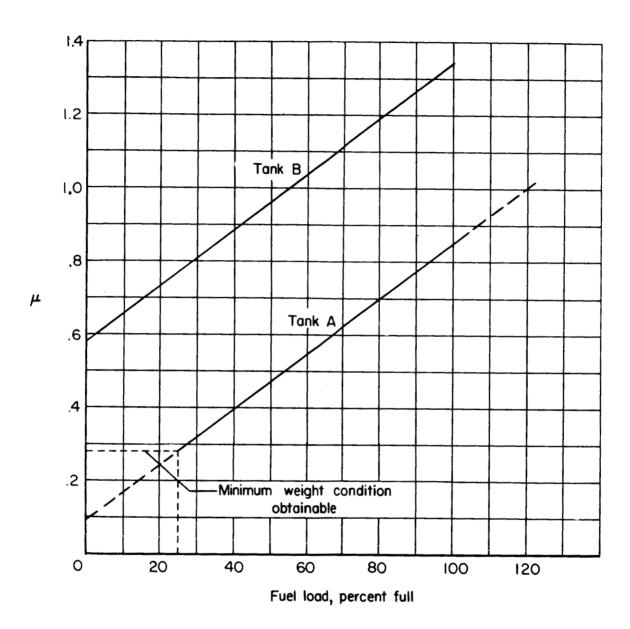
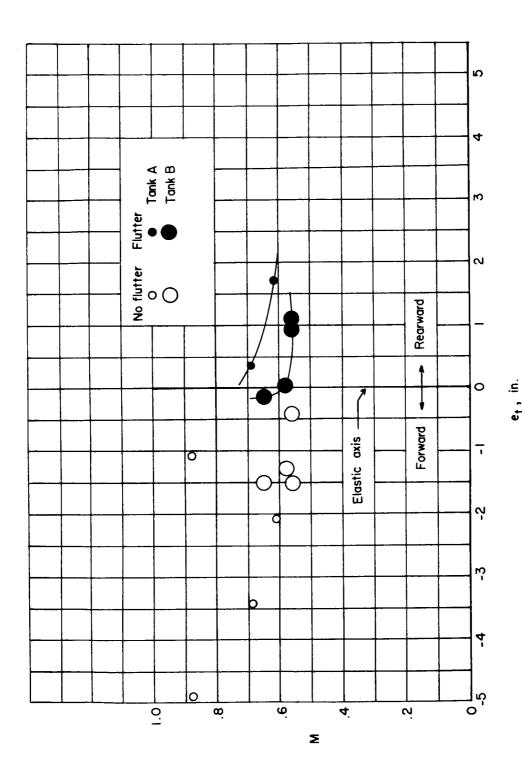


Figure 11.- Correlation of model tip-tank weight ratio with prototype fuel load.



gravity for both small and large tanks ballasted to the same weight and moment of inertia. $f_{h_1}/f_{\alpha_1}=0.92;~\mu=0.60;$ fuselage clamped. Figure 12.- Variation of flutter Mach number with tip-tank center of



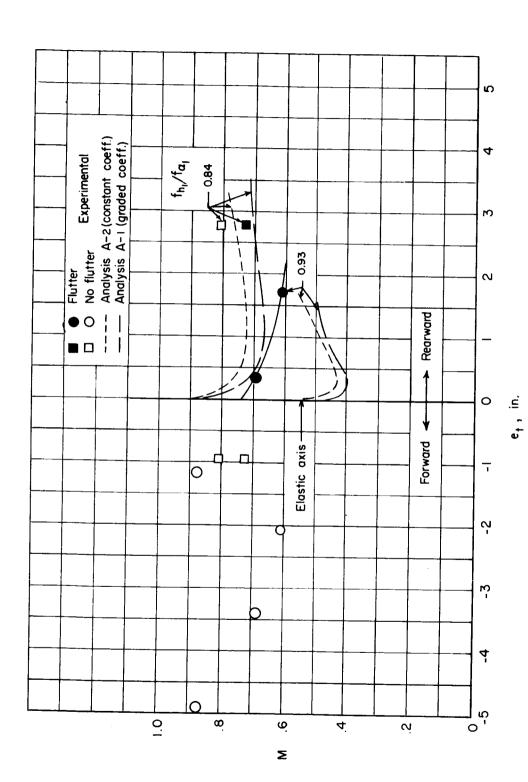
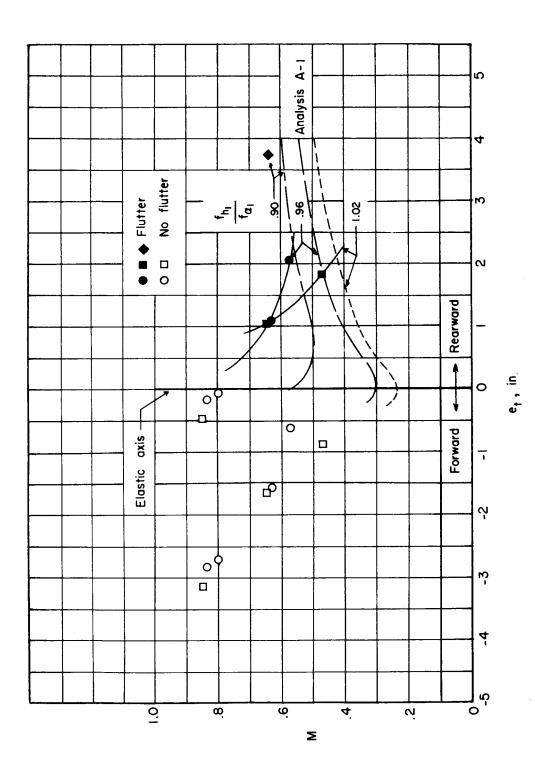


Figure 13.- Variation of experimental and theoretical flutter Mach numbers $f_{\rm hl}/f_{lpha_{
m l}}$. Tank A; with tip-tank center of gravity for two values of μ = 0.59; fuselage clamped.



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Figure 14.- Variation of experimental and theoretical flutter Mach numbers with tip-tank center of gravity for three values of $f_{\rm hl}/f_{\alpha_{\rm l}}$. Tank A; μ = 0.82; fuselage clamped.

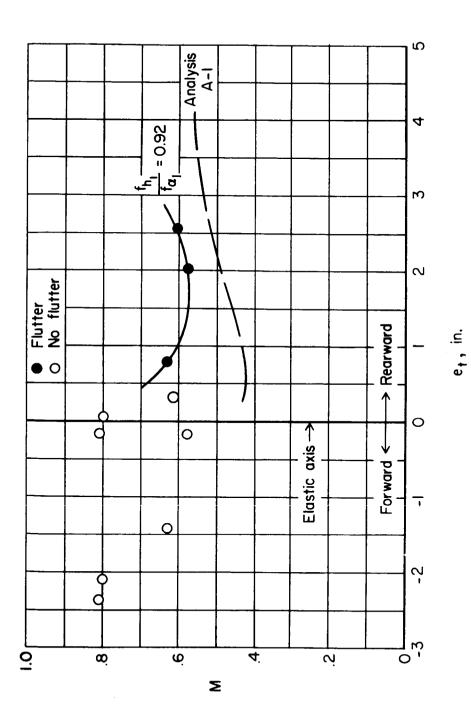


Figure 15.- Variation of experimental and theoretical flutter Mach numbers with tip-tank center of gravity. Tank A; μ = 1.005; fuselage clamped.

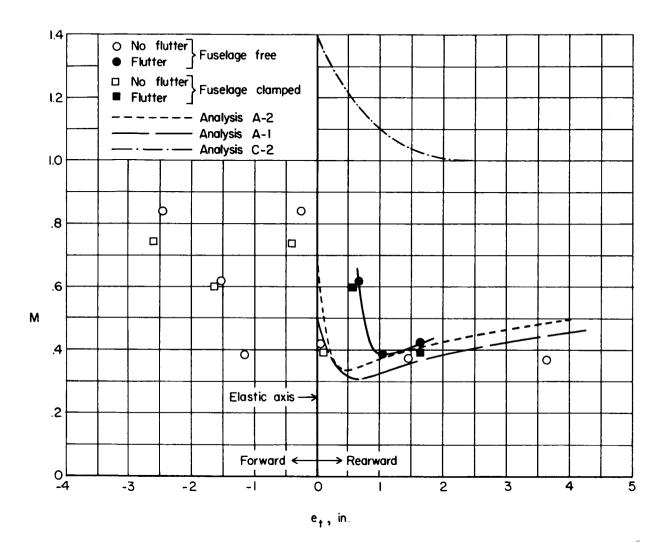


Figure 16.- Variation of experimental and theoretical flutter Mach numbers with tip-tank center of gravity with fuselage clamped and fuselage free to roll. Tank B; μ = 1.02; f_{h_1}/f_{α_1} = 1.04.

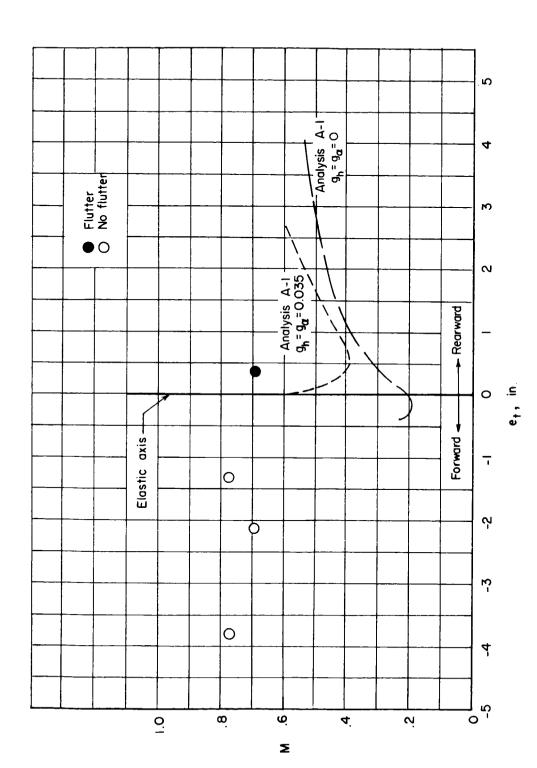


Figure 17.- Variation of experimental and theoretical Mach numbers with tip-tank center of gravity for $f_{h_1}/f_{\alpha_1}=0.98$. Tank B; $\mu=0.895$; fuselage clamped.

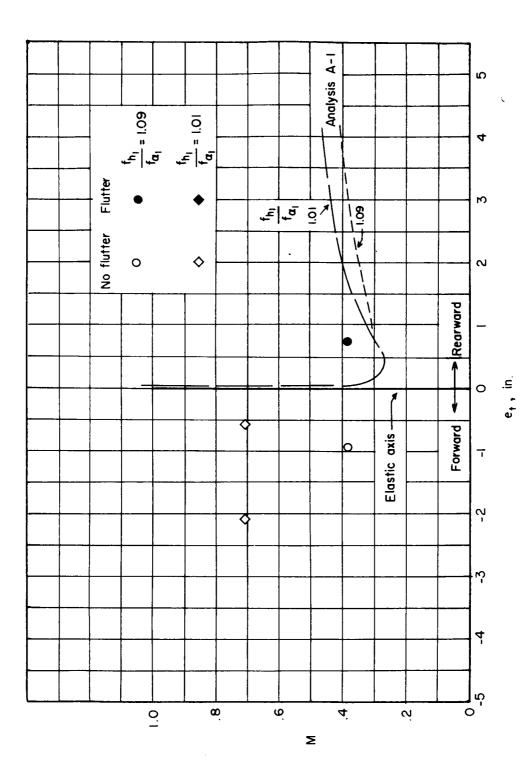


Figure 18.- Variation of experimental and theoretical flutter Mach numbers with tip-tank center of gravity. Tank B; $\mu = 1.29$; fuselage clamped.

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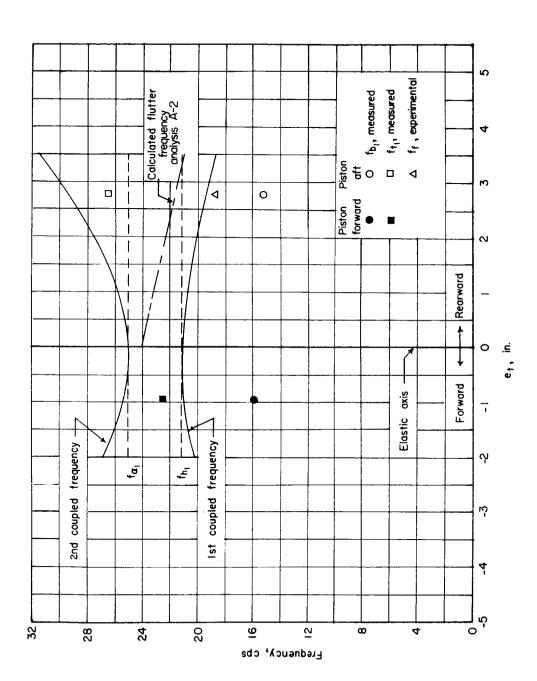


Figure 19.- Comparison of analytical and measured frequencies as a function of tip-tank center of gravity. Tank A; μ = 0.595; $f_{\rm hl}/f_{\alpha l}$ = 0.84; fuselage clamped.

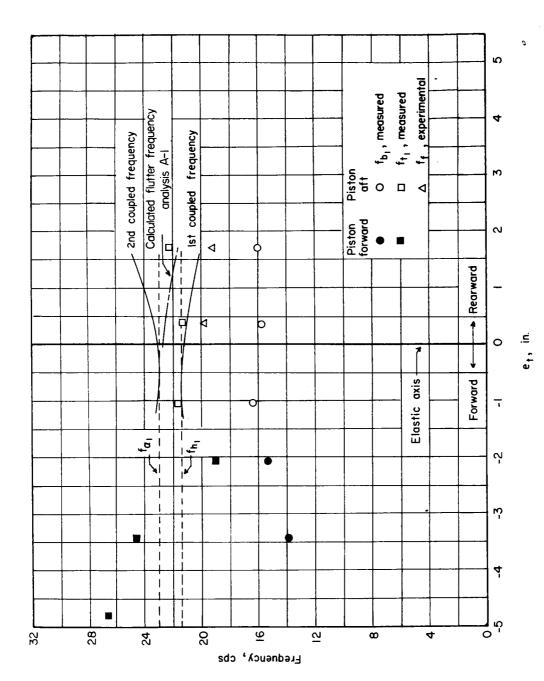


Figure 20.- Comparison of analytical and measured frequencies as a function of tip-tank center of gravity. Tank A; μ = 0.59; f_{h_1}/f_{α_1} = 0.93; fuselage clamped.

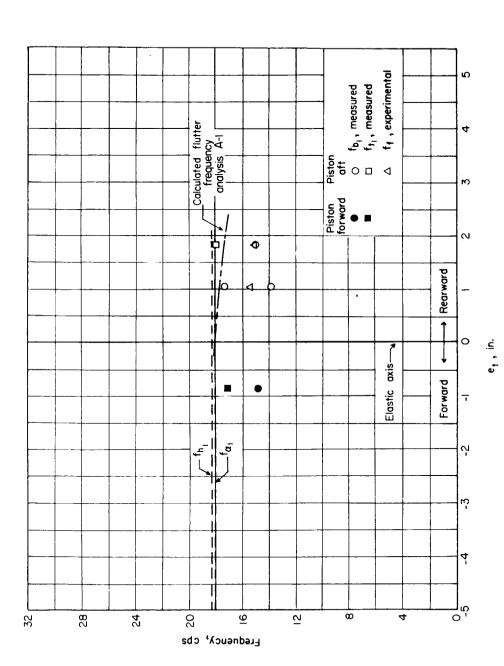


Figure 21.- Comparison of analytical and measured frequencies as a function of tip-tank center of gravity. Tank A; $\mu=0.82$; $f_{h1}/f_{\alpha_1}=1.02$; fuselage clamped.

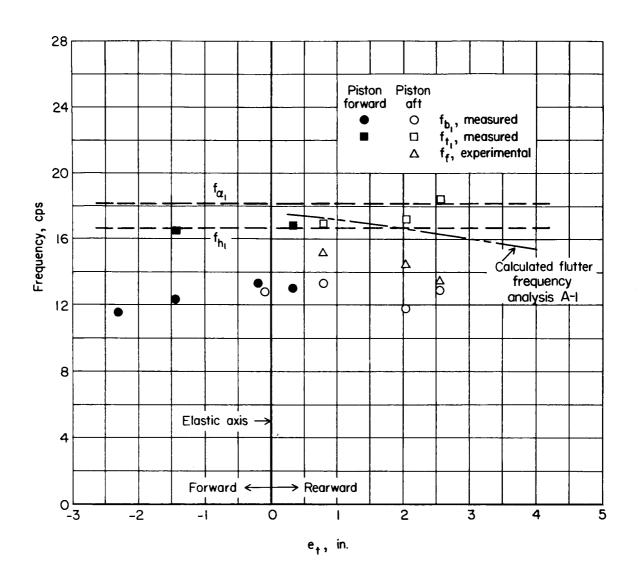


Figure 22.- Comparison of analytical and measured frequencies as a function of tip-tank center of gravity. Tank A; μ = 1.005; f_{h_1}/f_{α_1} = 0.92; fuselage clamped.

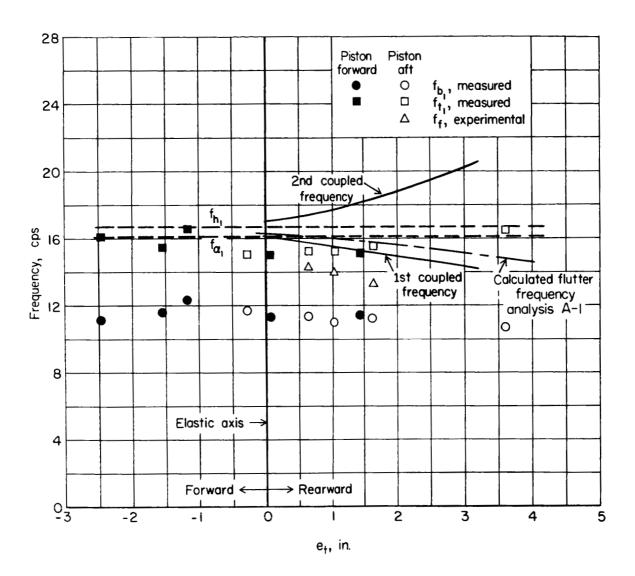


Figure 23.- Comparison of analytical and measured frequencies as a function of tip-tank center of gravity. Tank B; μ = 1.02; f_{h_1}/f_{α_1} = 1.04; fuselage clamped.

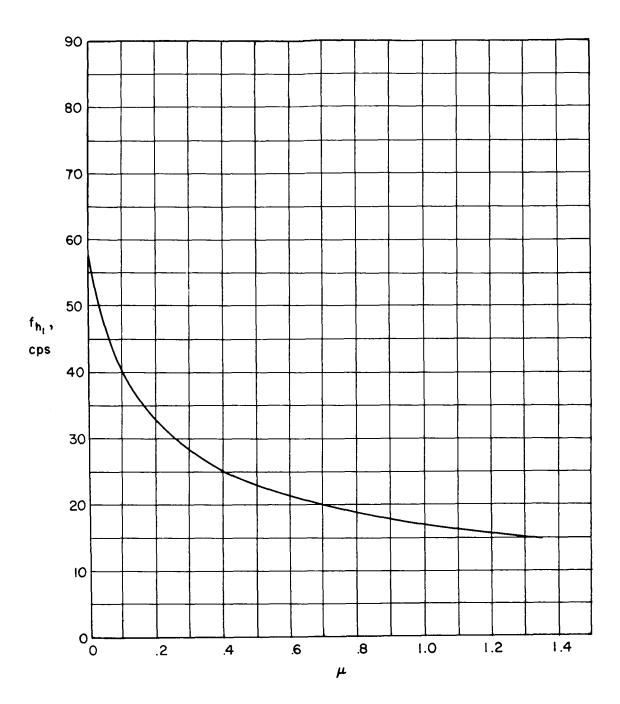


Figure 24.- Variation of first uncoupled bending frequency with tip-tank weight ratio.

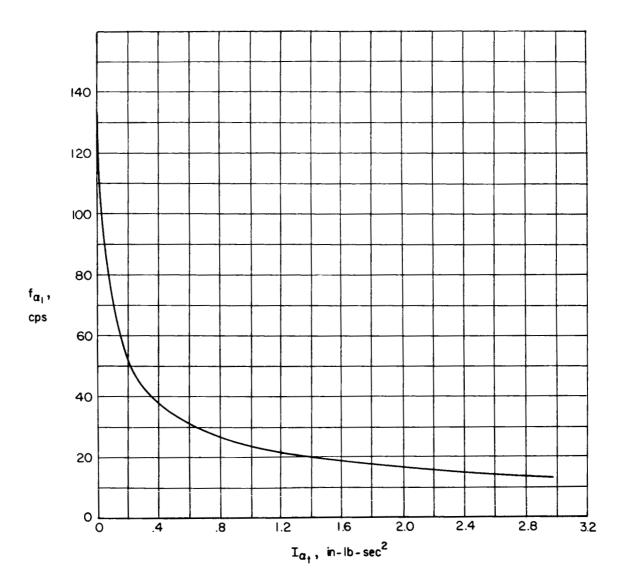


Figure 25.- Variation of first uncoupled torsional frequency with tiptank moment of inertia.

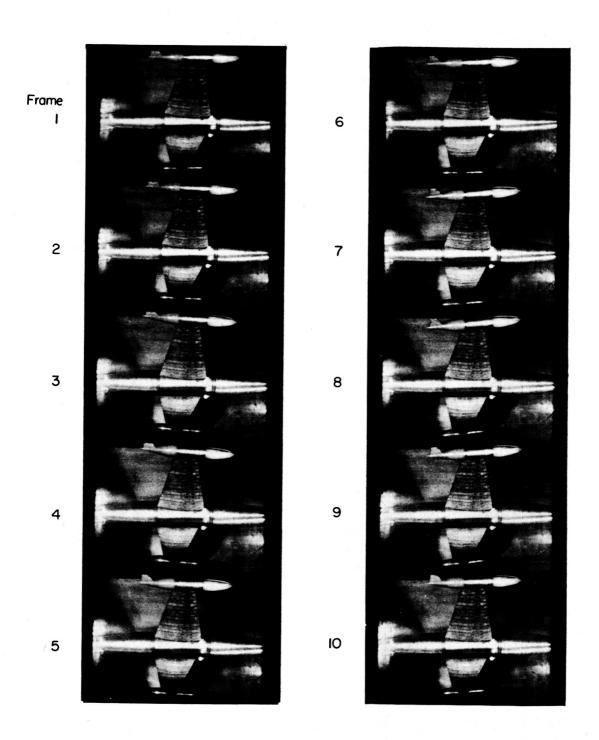


Figure 26.- High-speed motion pictures showing a typical cycle of symmetric flutter. The direction of increasing time is from top to bottom beginning with frame 1.

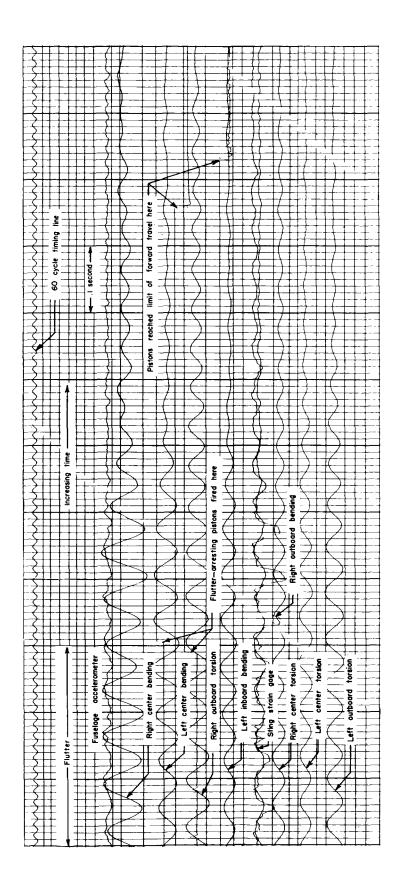


Figure 27.- Portion of oscillogram taken at flutter showing effect of flutter-arresting device.

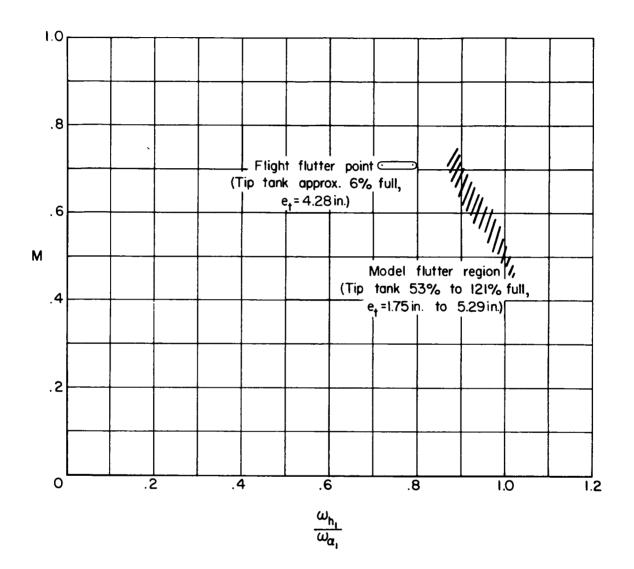


Figure 28.- Mach number against frequency ratio for model flutter region and flight flutter experience.

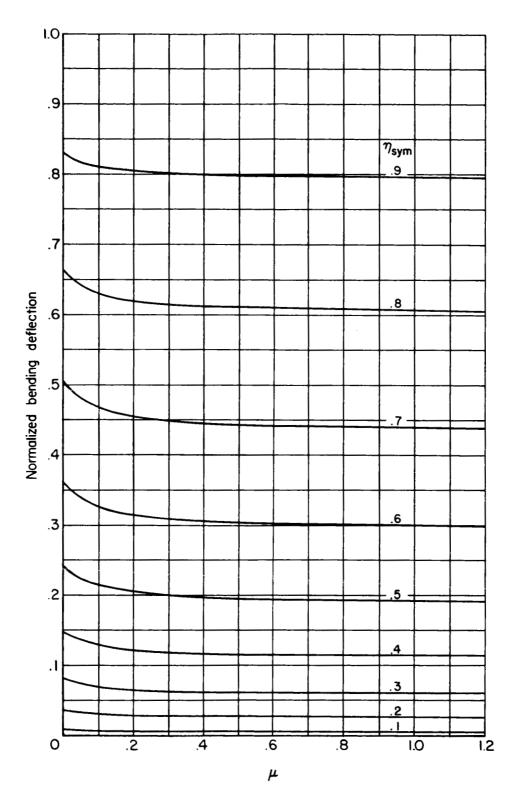


Figure 29.- Cross plot of uncoupled cantilever bending-mode shape as a function of tip-tank weight ratio.

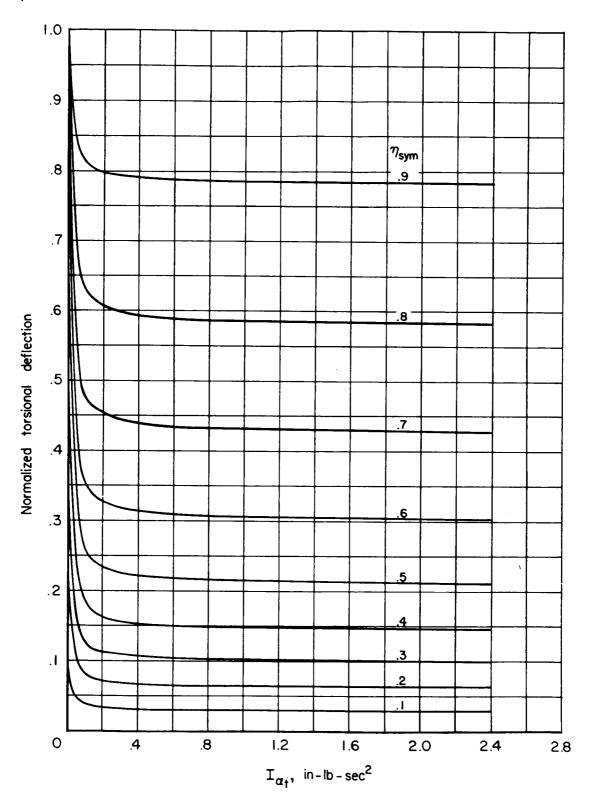


Figure 30.- Cross plot of uncoupled cantilever torsion-mode shape as a function of tip-tank moment of inertia.

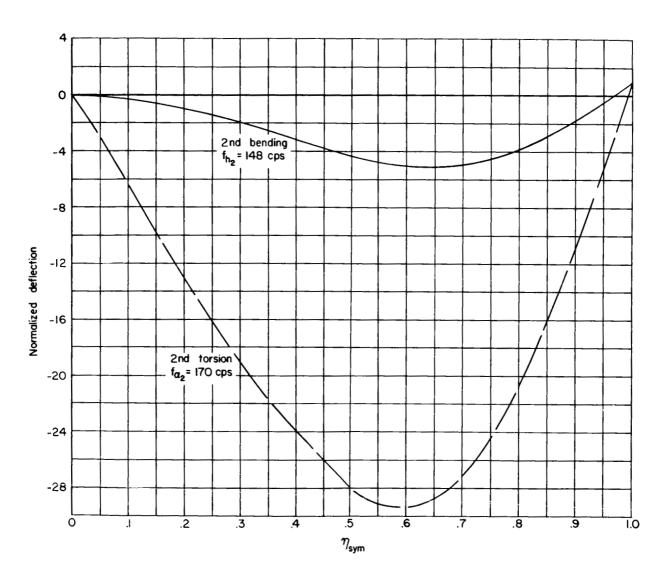
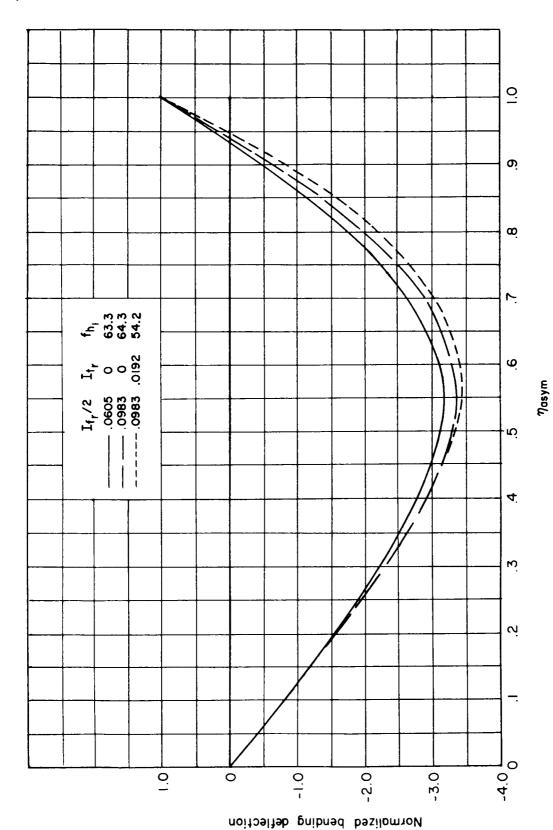
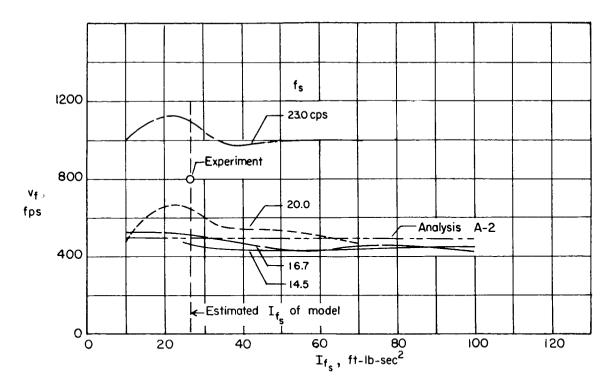


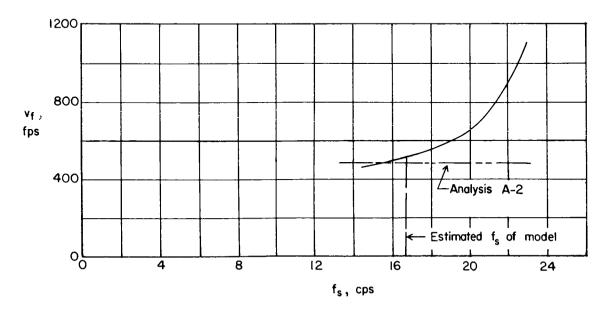
Figure 31.- Second uncoupled bending and torsion-mode shapes. Tank A; μ = 0.59; f_{h_1}/f_{α_1} = 0.93.



Tank B; Figure 32.- First uncoupled antisymmetric bending-mode shape. μ = 1.02.



(a) Effect of sting-fuselage pitching moment of inertia I_{f_S} on flutter speed for four values of sting-fuselage frequency f_S .



(b) Effect of sting-fuselage frequency on flutter speed for I_{fs} = 26.4.

Figure 33.- Correlation of experimental with theoretical flutter speeds including effects of sting flexibility (analysis B-1).

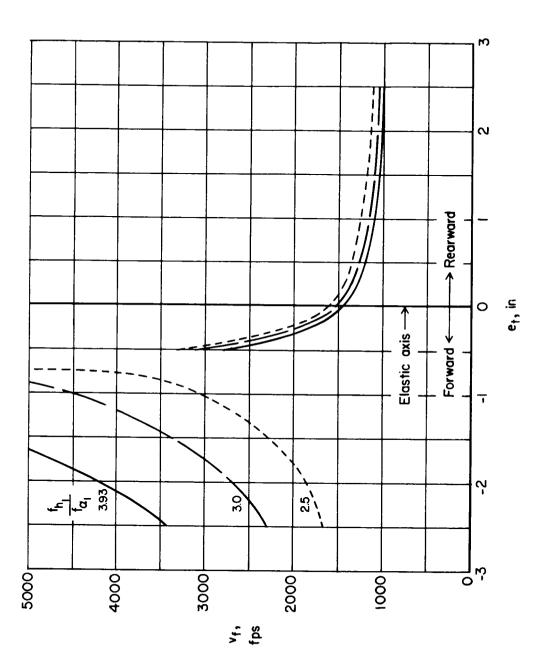


Figure 34.- Calculated antisymmetric flutter velocity (analysis C-1) against tip-tank center of gravity for various values of the frequency ratio. Tank B; μ = 1.02; If $_{\bf r}$ = 0.0605.

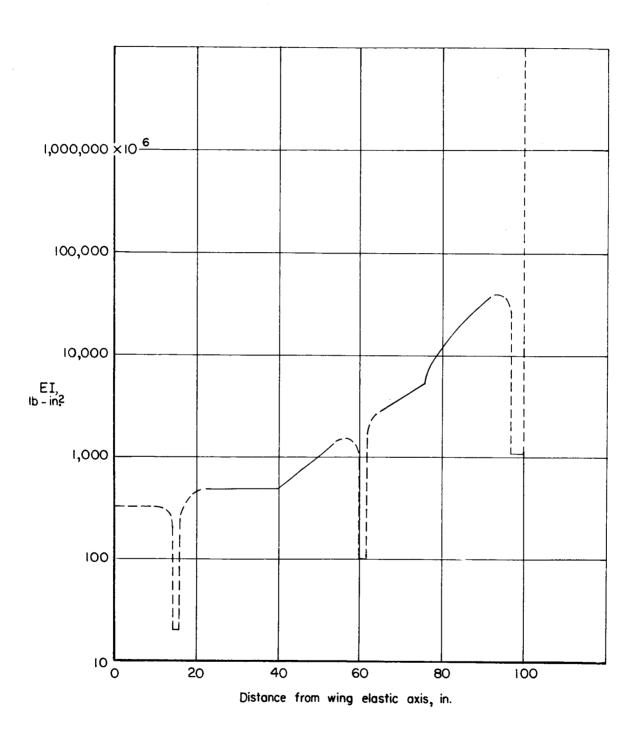


Figure 35.- Sting stiffness distribution.

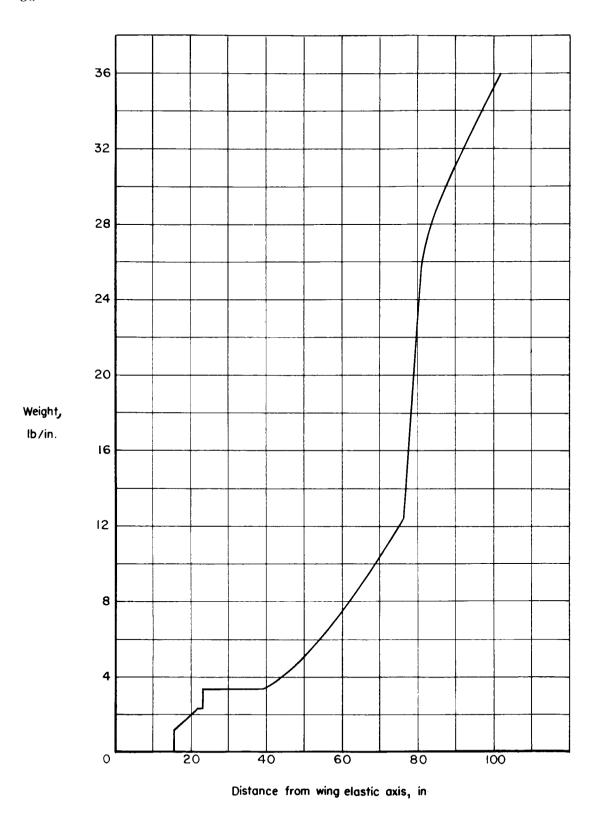


Figure 36.- Calculated sting weight distribution.